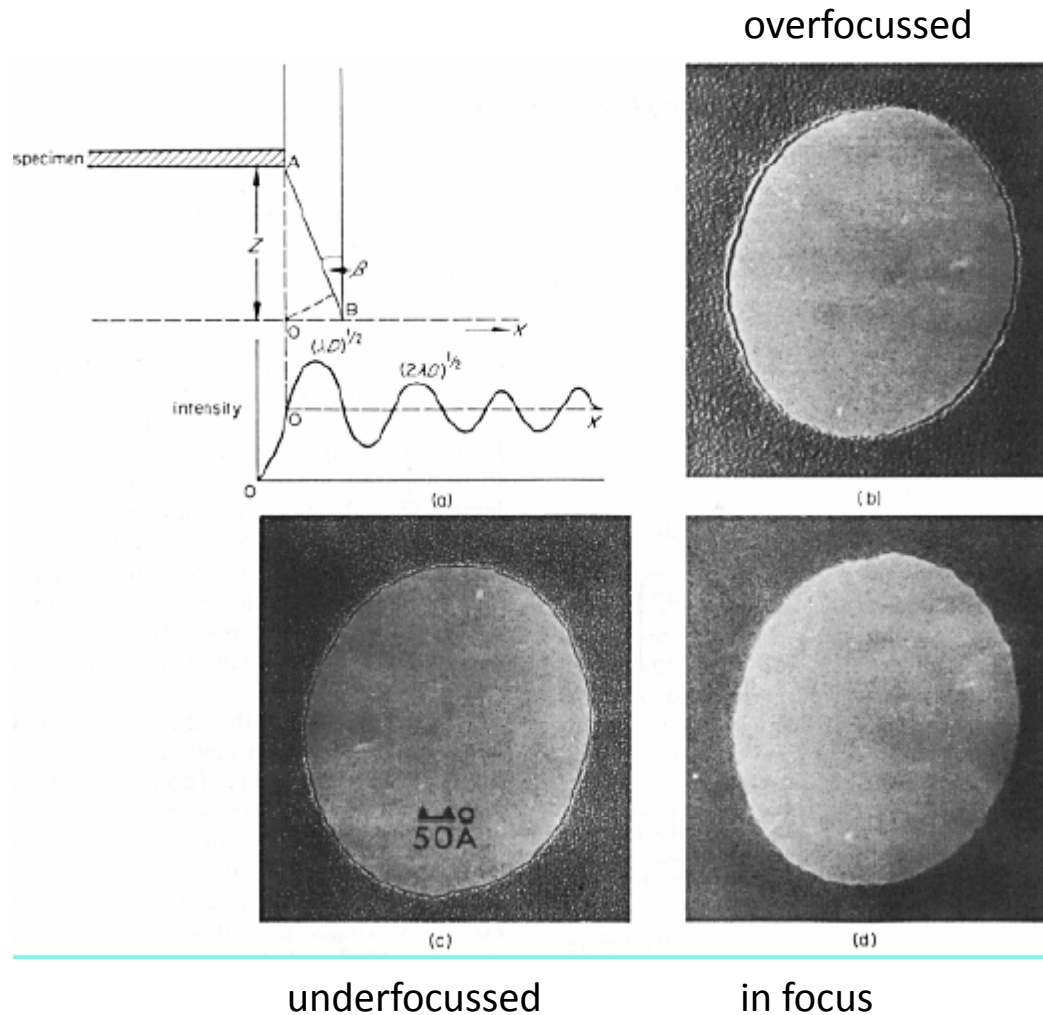


Phase microscopy: Fresnel fringes in the TEM

Interferences between waves not being in phase



2 waves propagating through different media (sample, vacuum)

=> difference of mean inner potentials

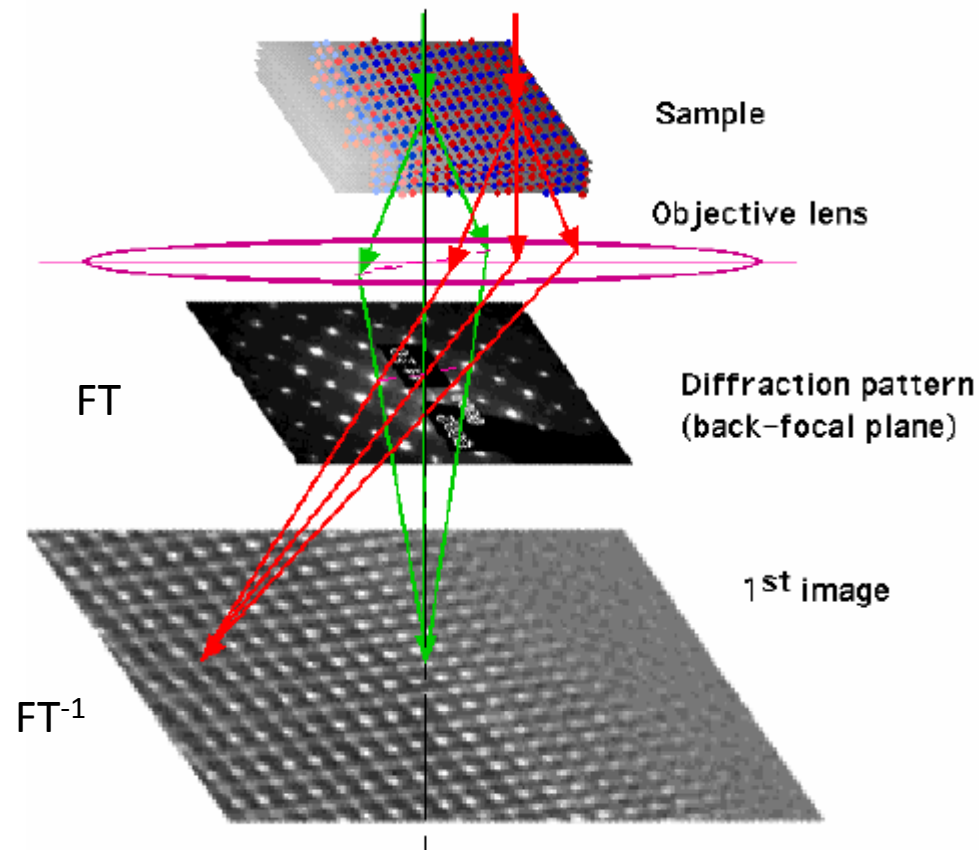
=> difference of optical paths

=> phase differences

Transformed into intensity contrast by the objective lens

Phase TEM relies on coherency of incident beam (energy spread, point source) => FEG

Phase microscopy: High Resolution Electron Microscopy



Multi-beam mode image: transmitted + several diffracted beams interfere
=> contrast = f (phase relationships of these beams)

- image “represents” the crystalline structure of the specimen => structural characterization at the atomic scale

- phase differences result from:
(i) interactions between the e-beam and the electrostatic potential in the object (periodical for a crystal)
(ii) phase shift induced by the objective lens (transfer function)

specimen thin enough => beam amplitude variations can be neglected (**“phase object”**)
=> phase shift proportional to the electrostatic potential at every point of the exit surface

High Resolution Electron Microscopy: WPO approx

- Transmission function of the sample:
if small phase shift, Weak Phase Object (WPO) approximation
 $q(x,y) = 1 - i\sigma\Phi(x,y)$ ← Projected potential

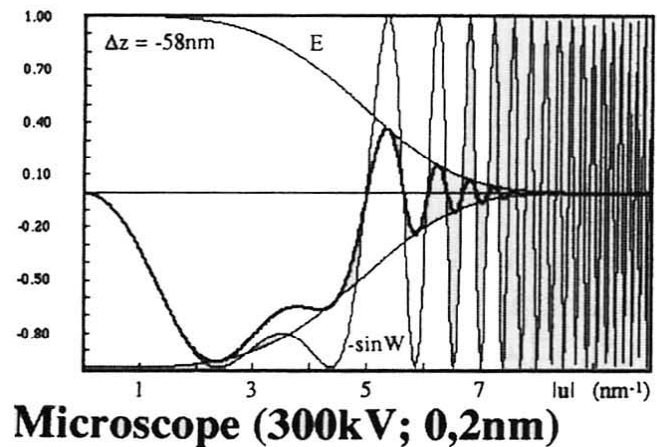
- Image directly related to the projected potential of the object (atoms) after modification by the objective lens (filter)

$$I(x,y) = |\psi_i(x,y)|^2 = |q(x,y) * t(x,y)|^2$$

$$= 1 + 2\sigma\Phi(x,y) * s(x,y)$$

Projected potential (object)
Im. of $t(x,y)$ (objective lens)

$s(x,y) = \text{TF}\{O(u,v)\sin(\chi(u,v))\}$ is the imaginary part of $t(x,y)$



- best “reading” of the projected potential when
 - largest possible frequency bandwidth
 - objective lens is underfocussed by $\Delta f = -1.2(Cs\lambda)^{1/2}$ (Scherzer condition)

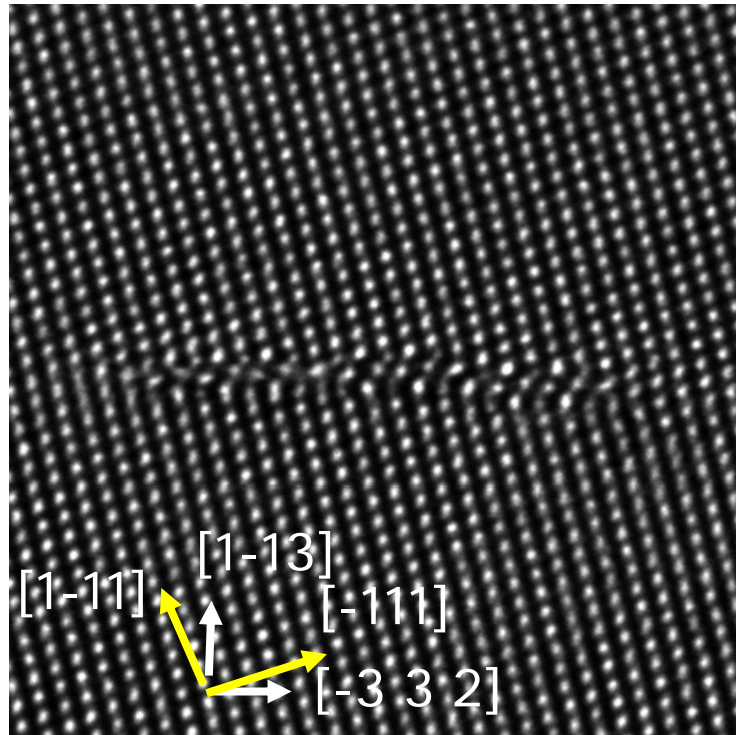
High Resolution Electron Microscopy

- In all practical cases, WPO approximation is not valid
=> comparison HREM images/numerical simulation of the images for $\neq D_f$ is necessary to access to the real structure of the sample (ex: where are the atomic columns?)
- However, although non linear, information regarding the atomic structure can be directly obtained in some cases: stacking sequences, extra plane, orientation relationship, presence of dislocations...
- Good experimental conditions:
 - thin specimen (10 nm range)
 - incident beam // to a simple (low index) direction of the crystal => diffraction= known plane of the RS
 - column perfectly aligned
 - appropriate defocus of the objective lens (Schertzer condition)
 - microscope resolution smaller than the distance between the atomic columns that you want to image!

Never trust ONE image !

High Resolution Electron Microscopy

Rod-like defect in silicon

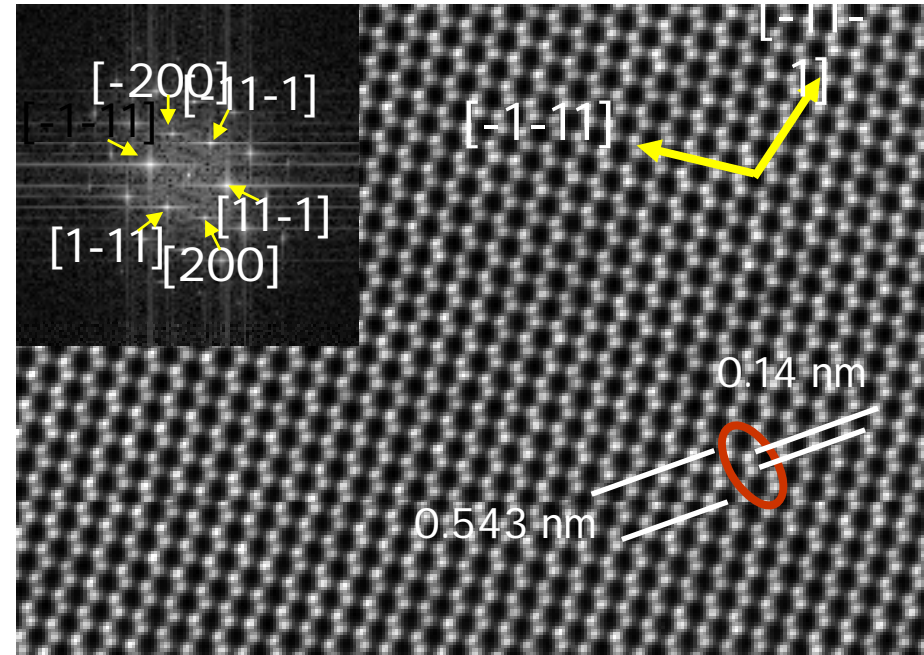


Cherkashin 2005

Interfaces, precipitates, defects, phase identification, defects, strain, etc...

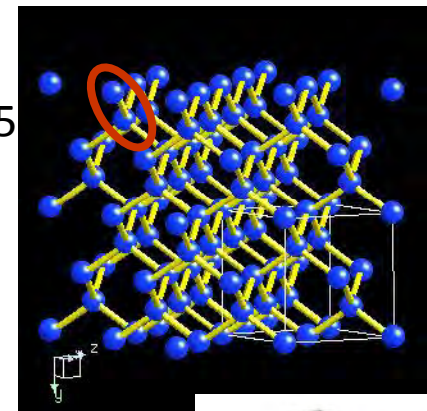
B=[011]

Si « dumbbells »

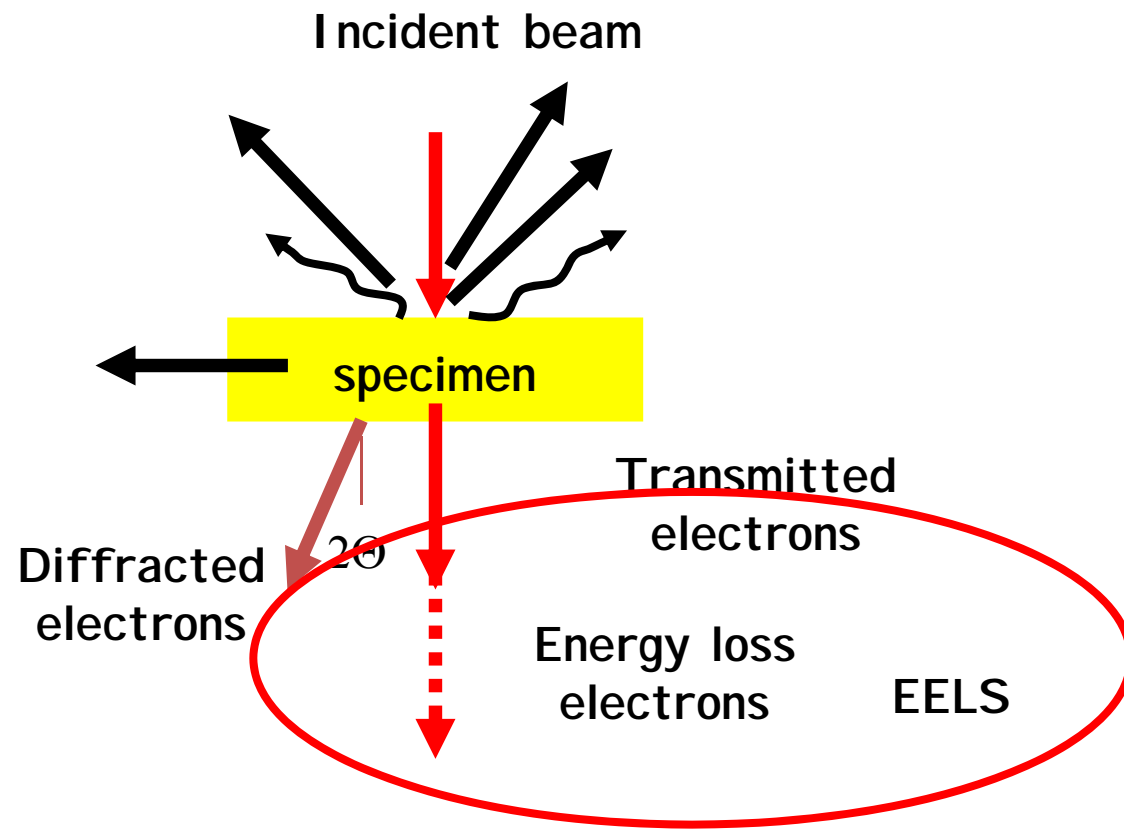


with Cs corrector

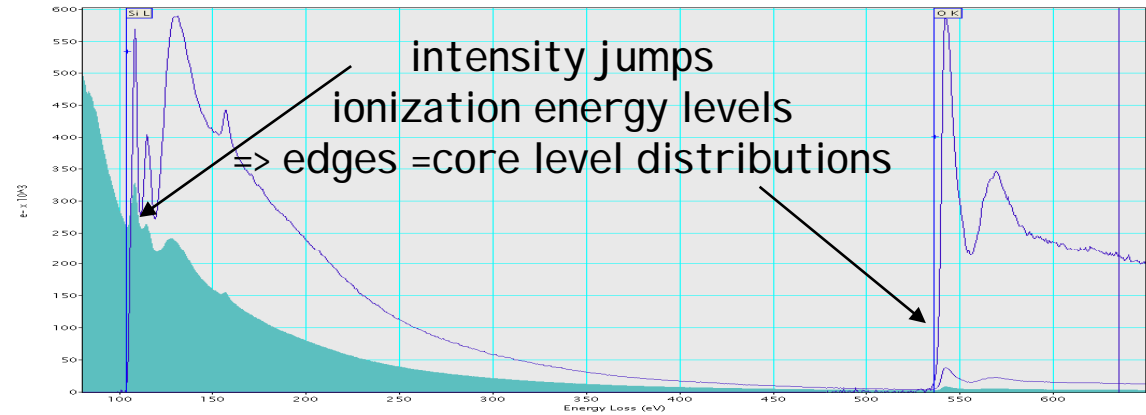
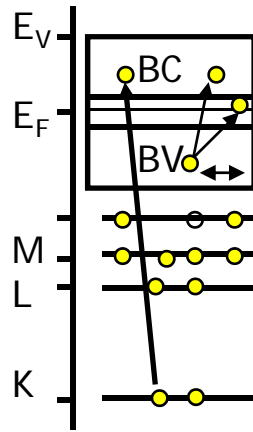
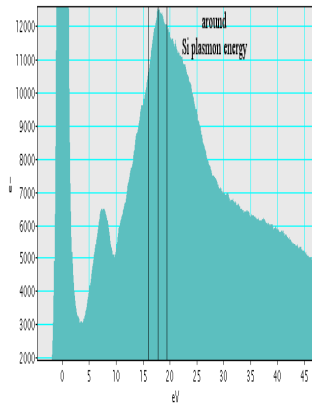
Houdellier 2005



What else ?



Electron Energy Loss Spectroscopy



$$\text{Im}(-1/\epsilon) = \epsilon_2 / (\epsilon_1^2 + \epsilon_2^2)$$

Low-losses (the first few ten eV)

Valence electrons (VEELS)

Collective excitation : Plasmon

Individual excitation : Interband Transition

Thickness: $I_t/I_{ZI} = t/\lambda_p$

Dielectric function

$$\epsilon = \epsilon_1 + i \epsilon_2$$

Optical parameters

$$n, k, \mu, R = f(\epsilon_1, \epsilon_2)$$

Electronic structure

$$\epsilon_2 \propto \text{JDOS}$$

Core-losses (up to a few thousand eV)

Inelastic interactions with inner or core level electrons

Elemental quantitative analysis

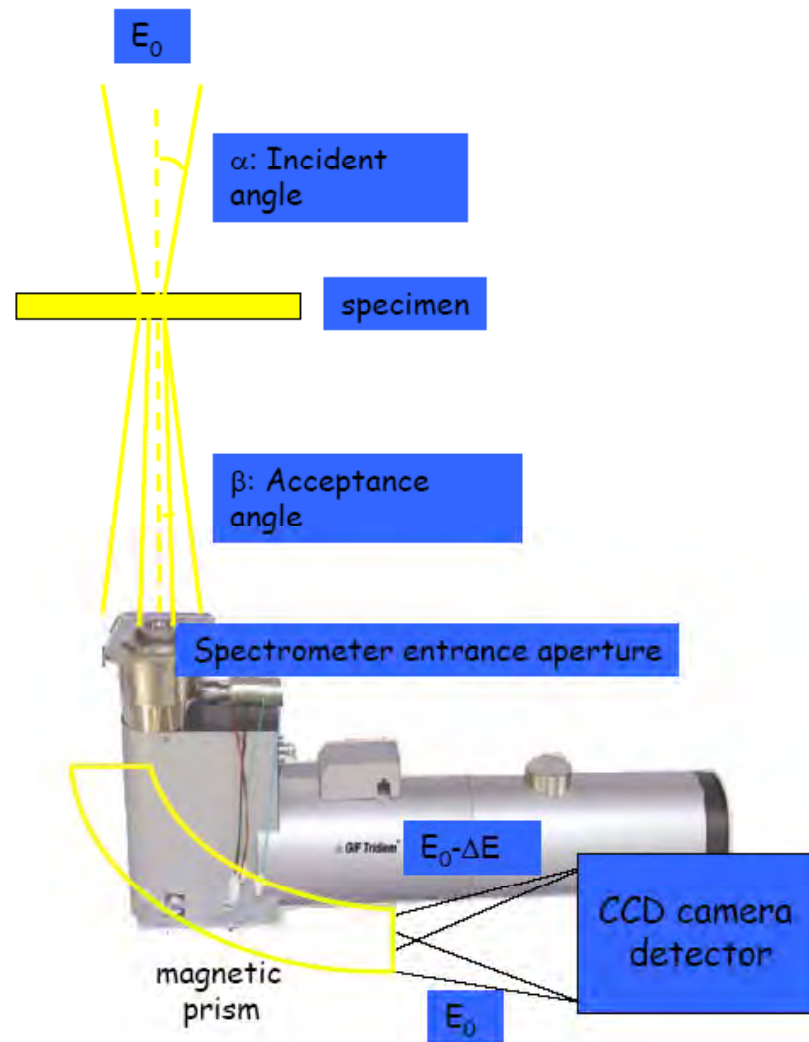
Electronic structure

$$I \sim \text{Im}(-1/\epsilon) = \epsilon_2 \propto \text{unoccupied DOS}$$

Energy Loss Near Edge fine Structure

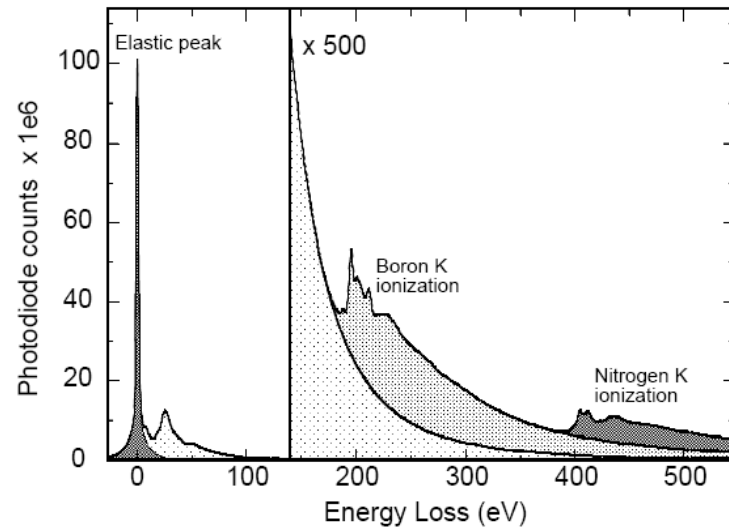
ELNES

Electron Energy Loss Spectroscopy

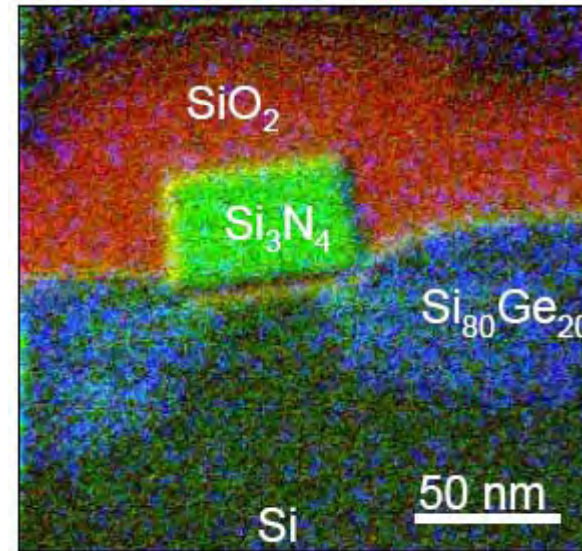


Electron Energy Loss Spectroscopy

. Electron Energy Loss Spectroscopy



Energy filtered TEM



TEM/EELS

EELS adds one dimension to TEM imaging

By probing the excitations of electrons bounded to the solid with the electrons from the incident beam

Gives information on:

- specimen thickness
- elemental chemical composition and spatial distribution of elements
- spatial distribution of first neighbors relative to an atomic site
- chemical bonding
- electronic band structure
- dielectric function

with an energy resolution ranging between 1 eV and 0.3 eV

with a spatial resolution ranging between 10 nm and 1 nm

Concept of holography

Complex wave : amplitude and phase

$$\Psi = Ae^{i\varphi}$$

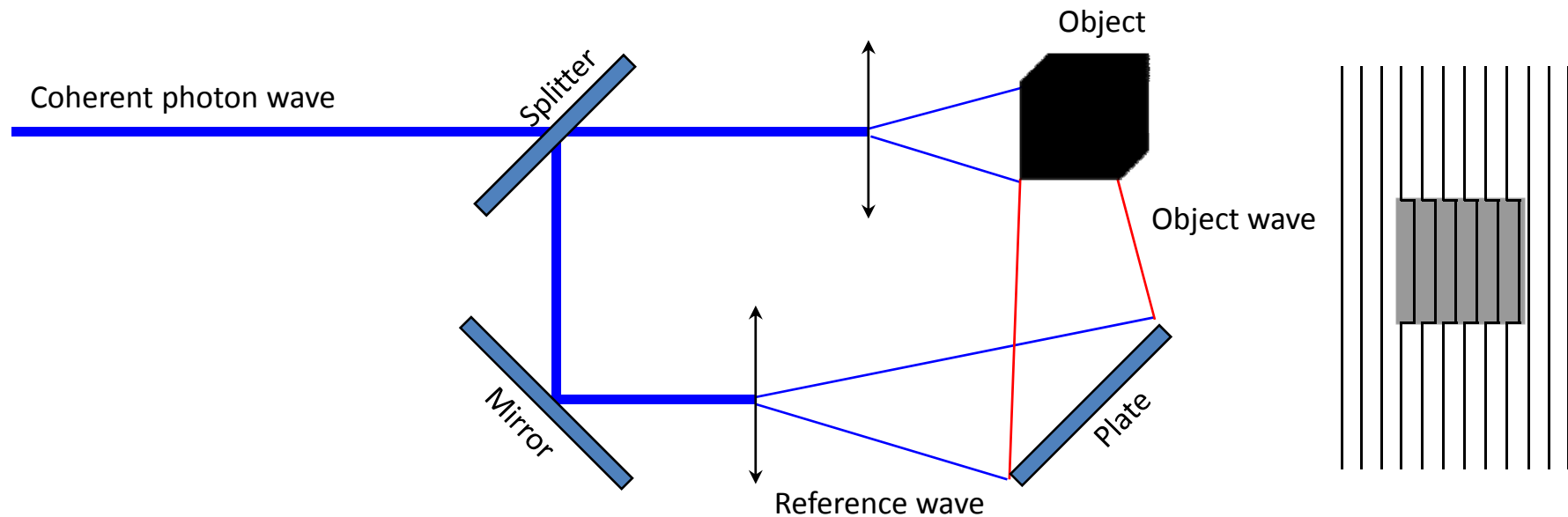
Take a picture :



$$I = |\Psi|^2 = A^2$$

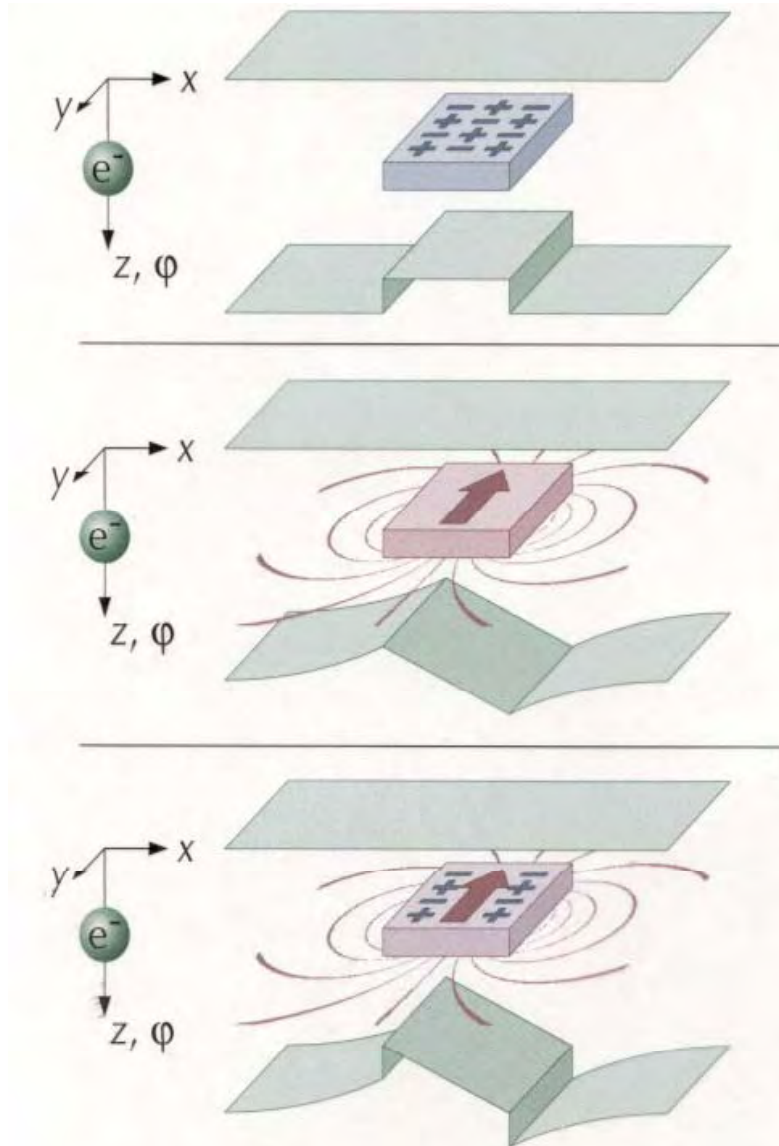
You loose the phase φ !!!

We want to retrieve the phase !



$$I = |\Psi_{ref} + \Psi_{object}|^2 = A_{ref}^2 + A_{object}^2 + 2A_{ref}A_{object} \cos(2\pi R_0 x + \varphi_{object} - \varphi_{ref})$$

Information in the phase

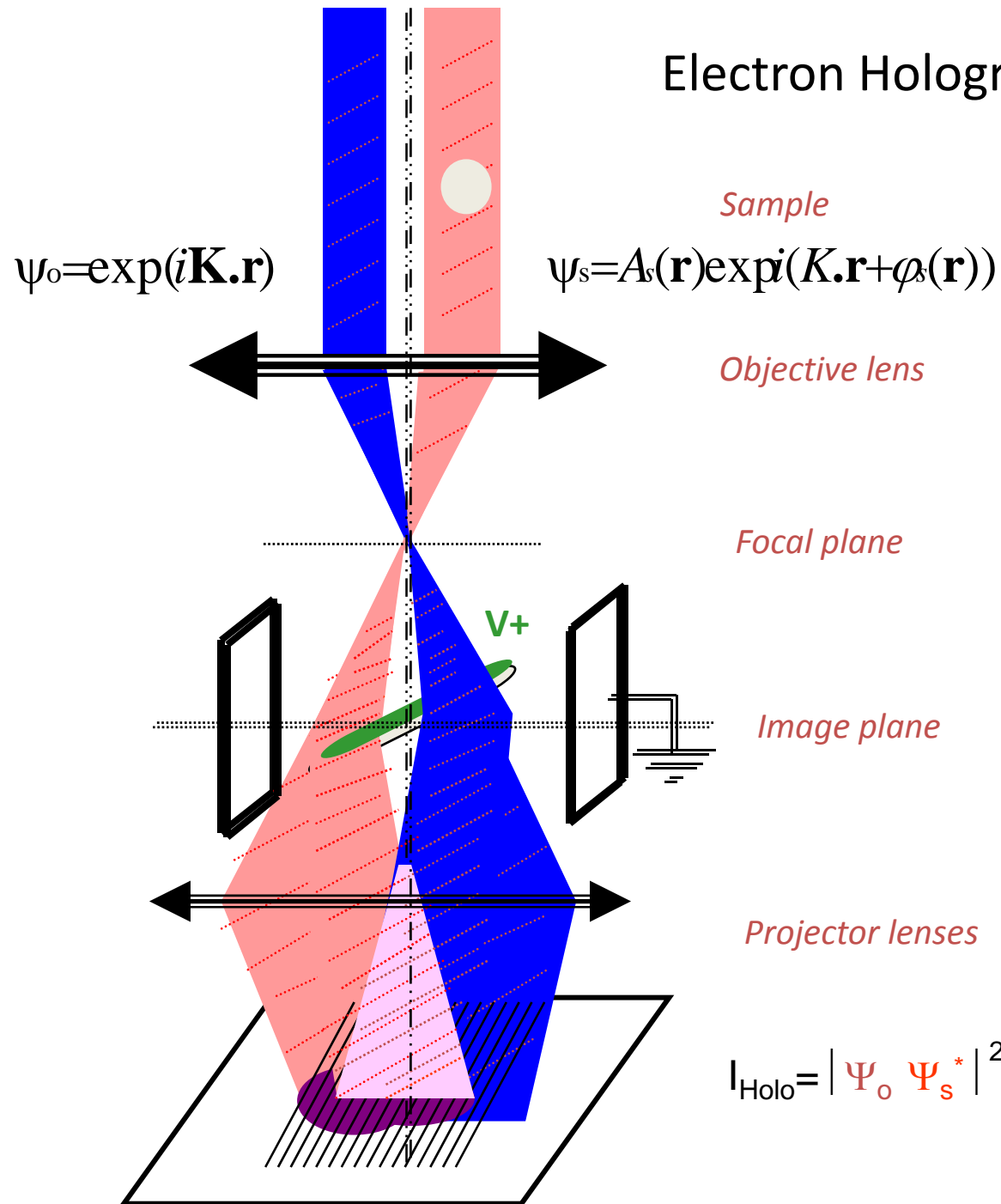


$$\phi^E(\mathbf{x}) = C_E \int V(x, z) dz$$

$$\phi^M(\mathbf{x}) = -\frac{e}{\hbar} \iint B_n(\mathbf{x}, z) dx dz$$

$$\phi^T = \phi^M + \phi^E$$

Electron Holography in a TEM



$$I_{\text{Holo}} = |\Psi_o \Psi_s^*|^2 = 1 + A_s^2(x, y)$$

$$+ 2A_s(x, y) \cos[2\pi \mathbf{R}_0 \cdot \mathbf{x} + \phi_s(x, y)]$$

GPA analysis of the hologram

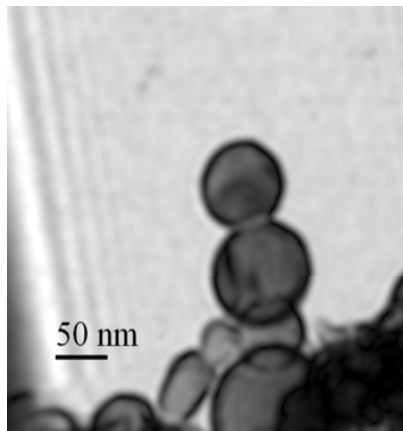
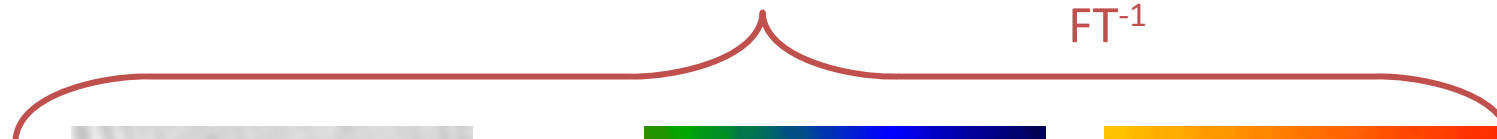
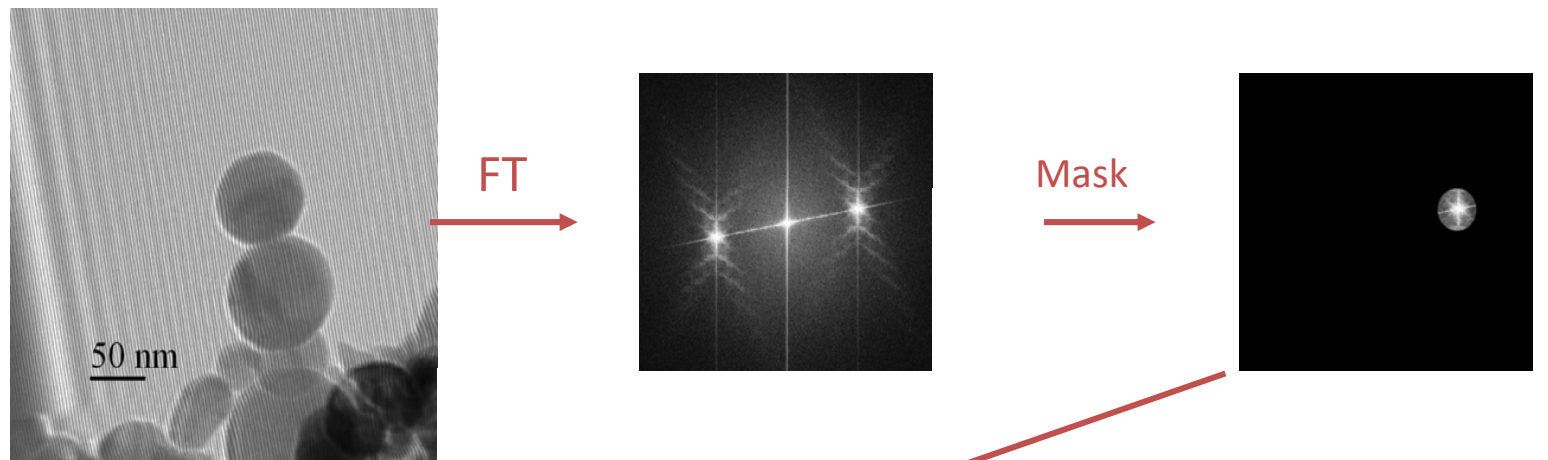
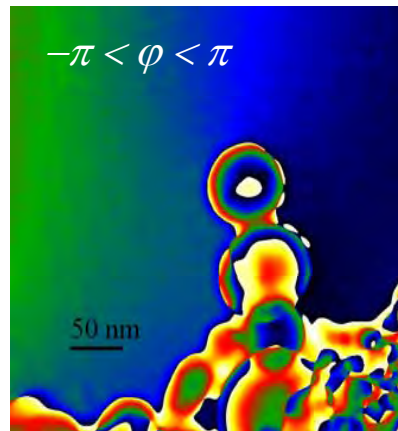
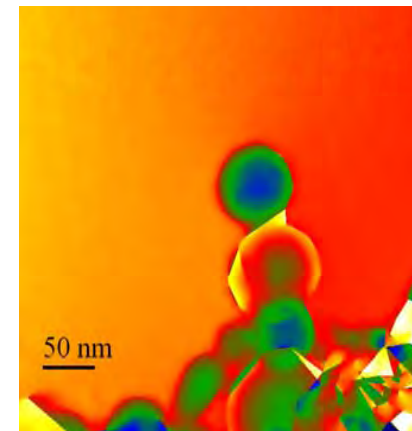


Image d'amplitude
 $A_S(x, y)$



2π jumps



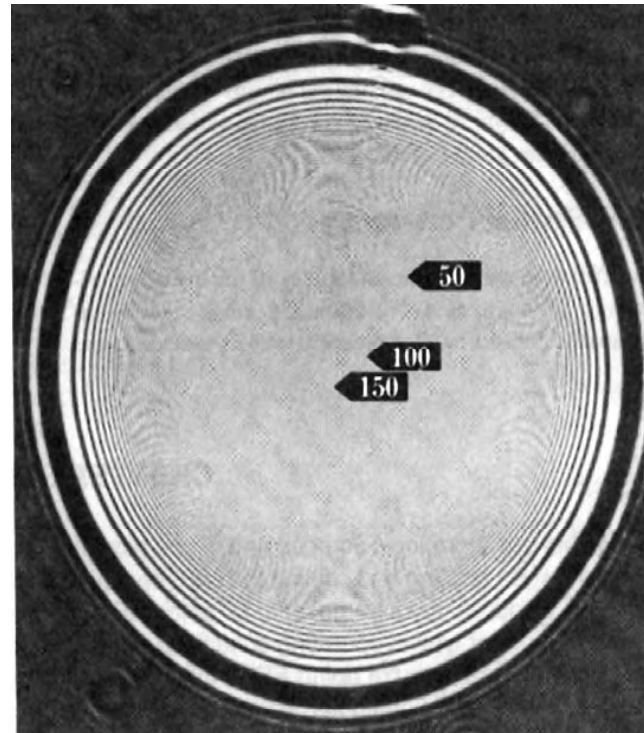
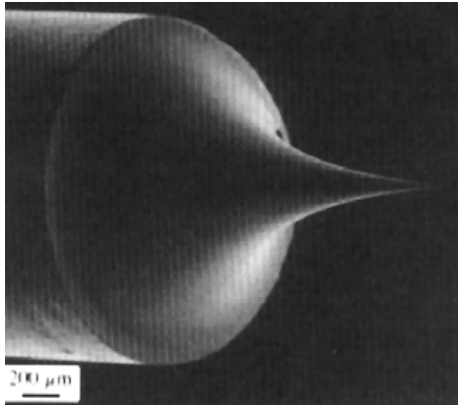
unwrapped

Image de phase
 $\varphi_S(x, y)$

Experimental aspects

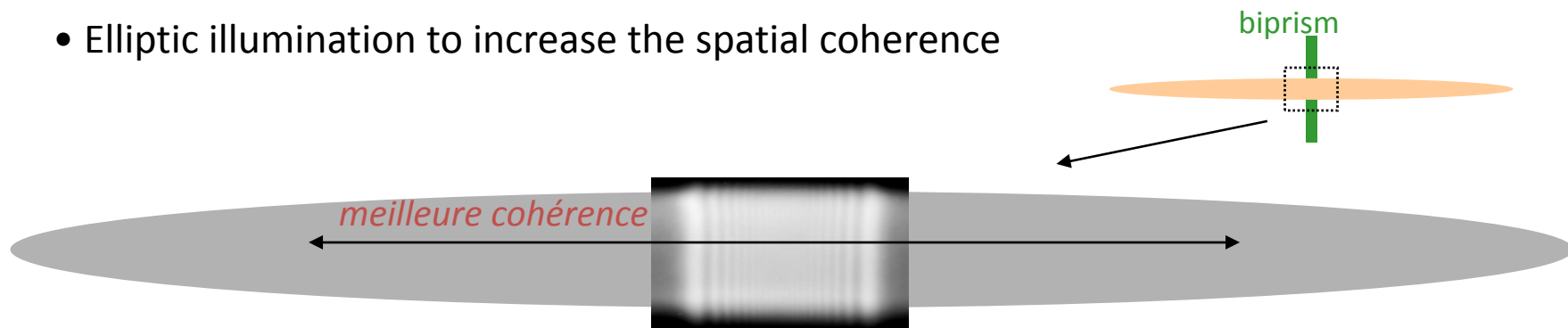
- Coherent Electron Beam

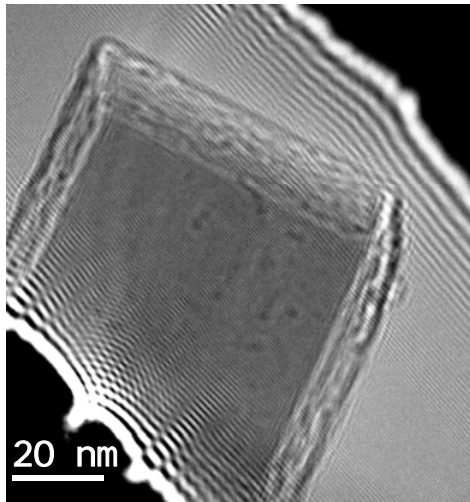
Field Emission Gun



Fresnel fringes
due to the high
coherence of
the beam

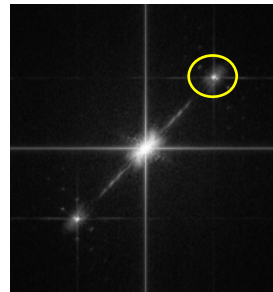
- Elliptic illumination to increase the spatial coherence



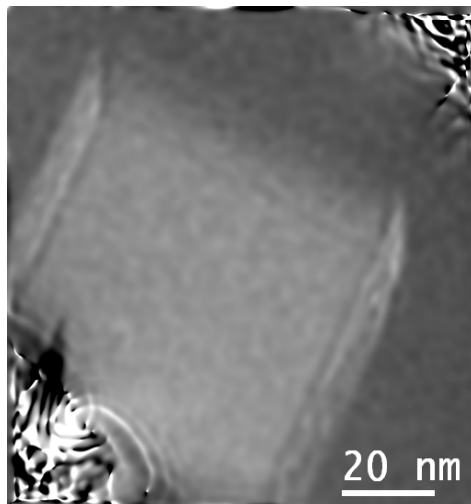


Hologram with the object

FT

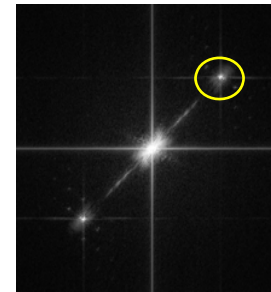


FT⁻¹

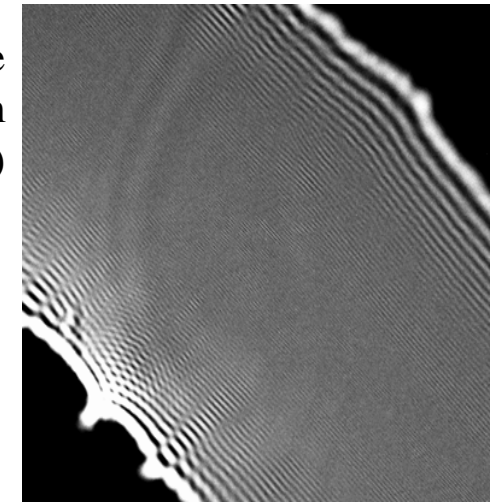


Reference
hologram
(without object)

FT



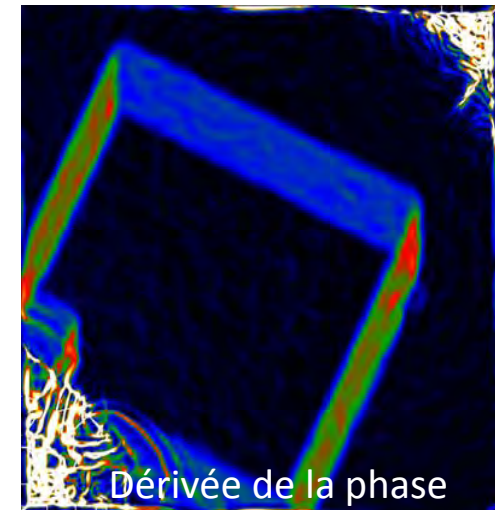
FT⁻¹



$$I_S = A_S e^{i\varphi_S}$$

$$I_R = A_R e^{i\varphi_R}$$

$$\frac{I_S}{I_R} = \frac{A_S}{A_R} e^{i(\varphi_S - \varphi_R)}$$



$$\text{Normalised Amplitude } \frac{A_S}{A_R} = A_N$$

$$\text{Phase shift due to the object only } \varphi_S - \varphi_R = \Delta\varphi$$

Non magnetic sample ($B = 0$)

$$\varphi(x, y) = C_E \int V(x, y, z) dz - \frac{e}{\hbar} \iiint B_n(x, y, z) dx dz$$

$$\varphi_{elect}(x, y) = C_E \int V(x, y, z) dz$$

$$\text{with : } C_E = \left(\frac{2\pi}{\lambda} \right) \left(\frac{E + E_0}{E(E + 2E_0)} \right)$$

$$V(x, y, z) = V_i(x, y, z) + V(x, y, z)$$

- $V_i(x, y, z)$ = Mean Inner Potential (MIP)
- $V(x, y, z)$ = Local electric potential (junctions, CMOS, diode ..)
- if V_i et V are constants along the e-beam direction « z » :

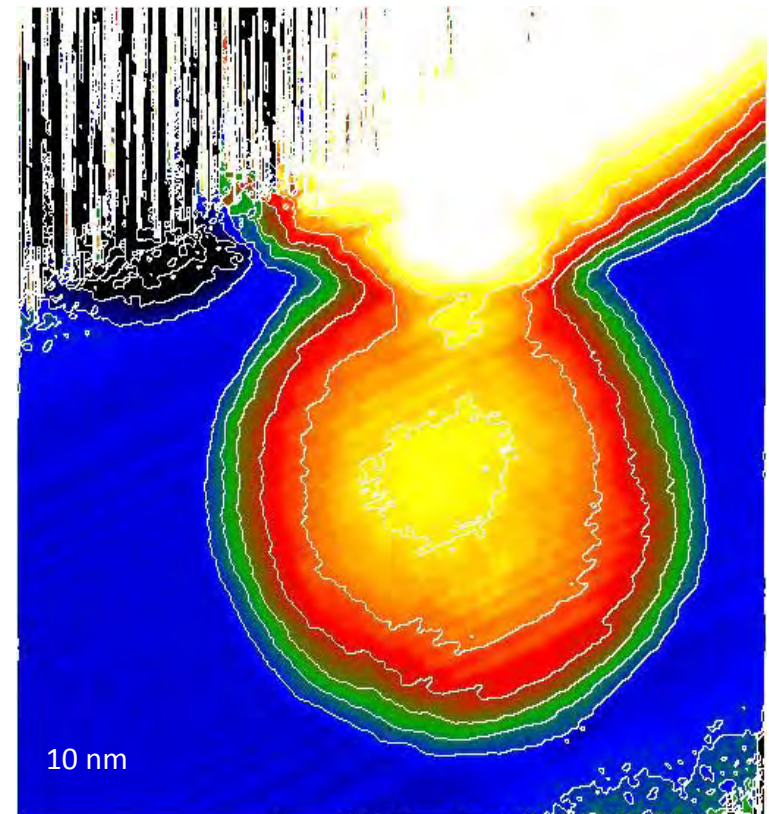
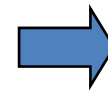
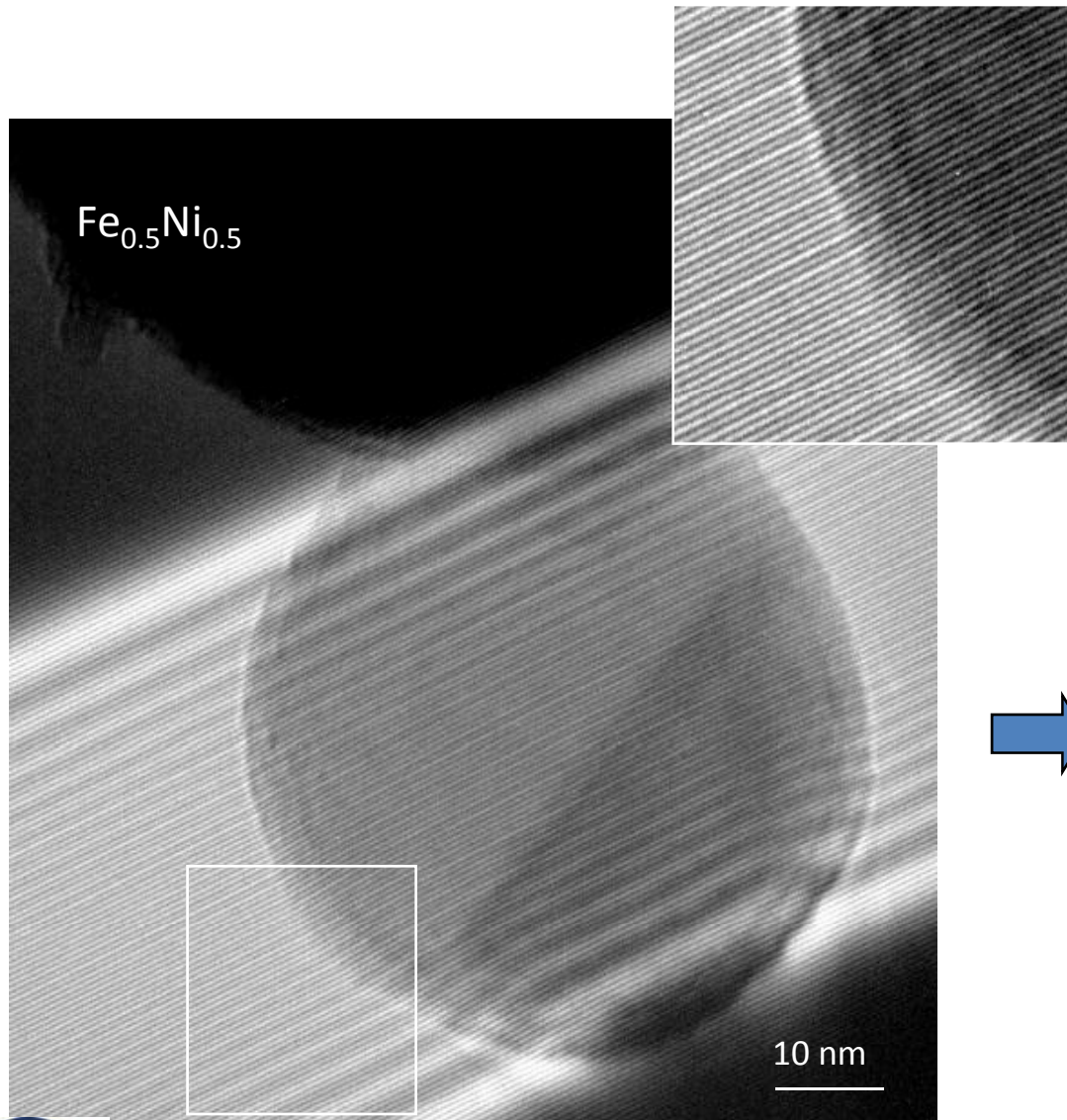
$$\varphi_{elect} = C_E (V_i + V) t$$

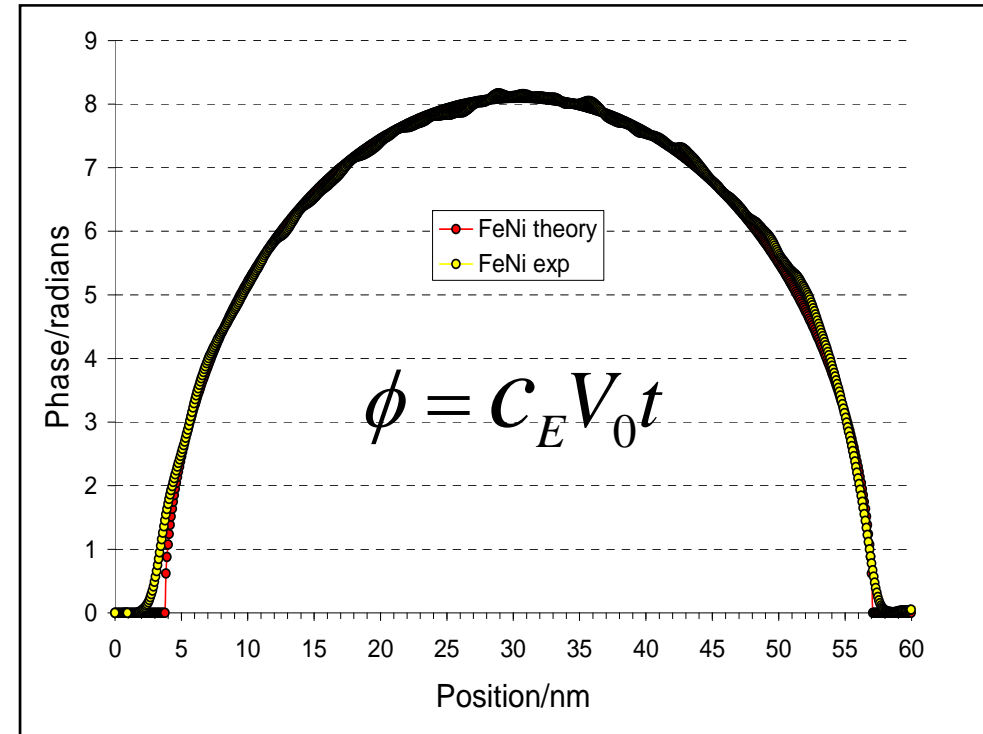
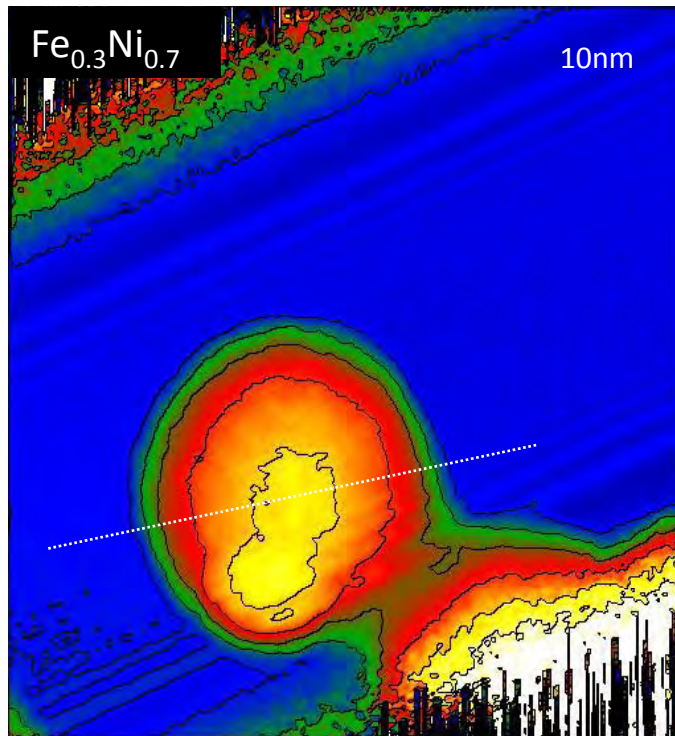
t = sample thickness

➔ Need to perfectly know the sample thickness (sphere, cleaved sample, CBED thickness measurement...) to quantify V_i and V

Mean Inner Potential (MIP) measurement

$$\varphi_{\text{elect}} = C_E V_i t$$



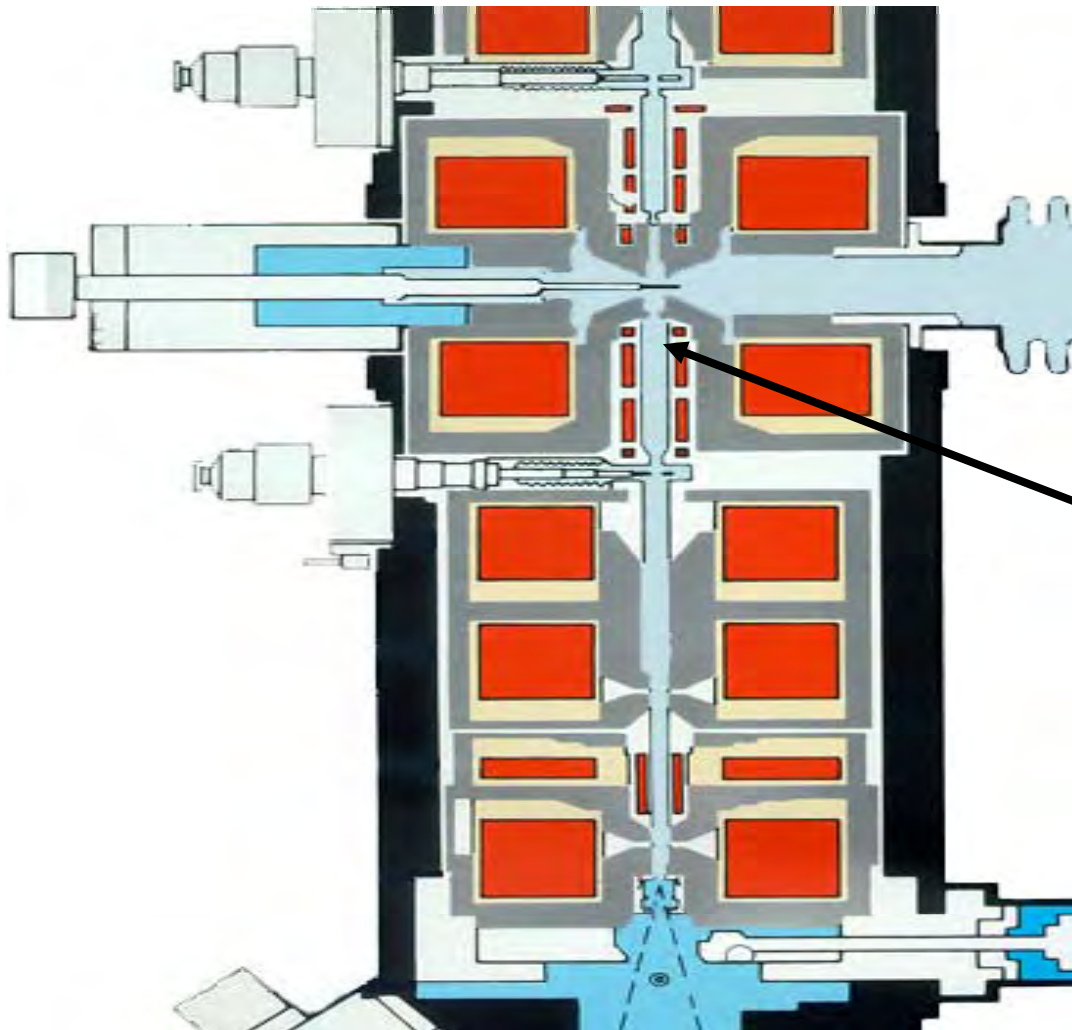


- MIP measurement
- Composition analysis
- Morphological studies

Magnetic samples

$$\varphi(\mathbf{x}) = C_E \int V(\mathbf{x}, z) dz - \frac{e}{\hbar} \iint B_n(\mathbf{x}, z) dx dz$$

Pb. 1: The objective must be switched off. => use of Lorentz lens to get sufficient resolution



- Object in a magnetic field free area

- Field of view up to $1\mu\text{m}$

$f = 30 \text{ mm}$

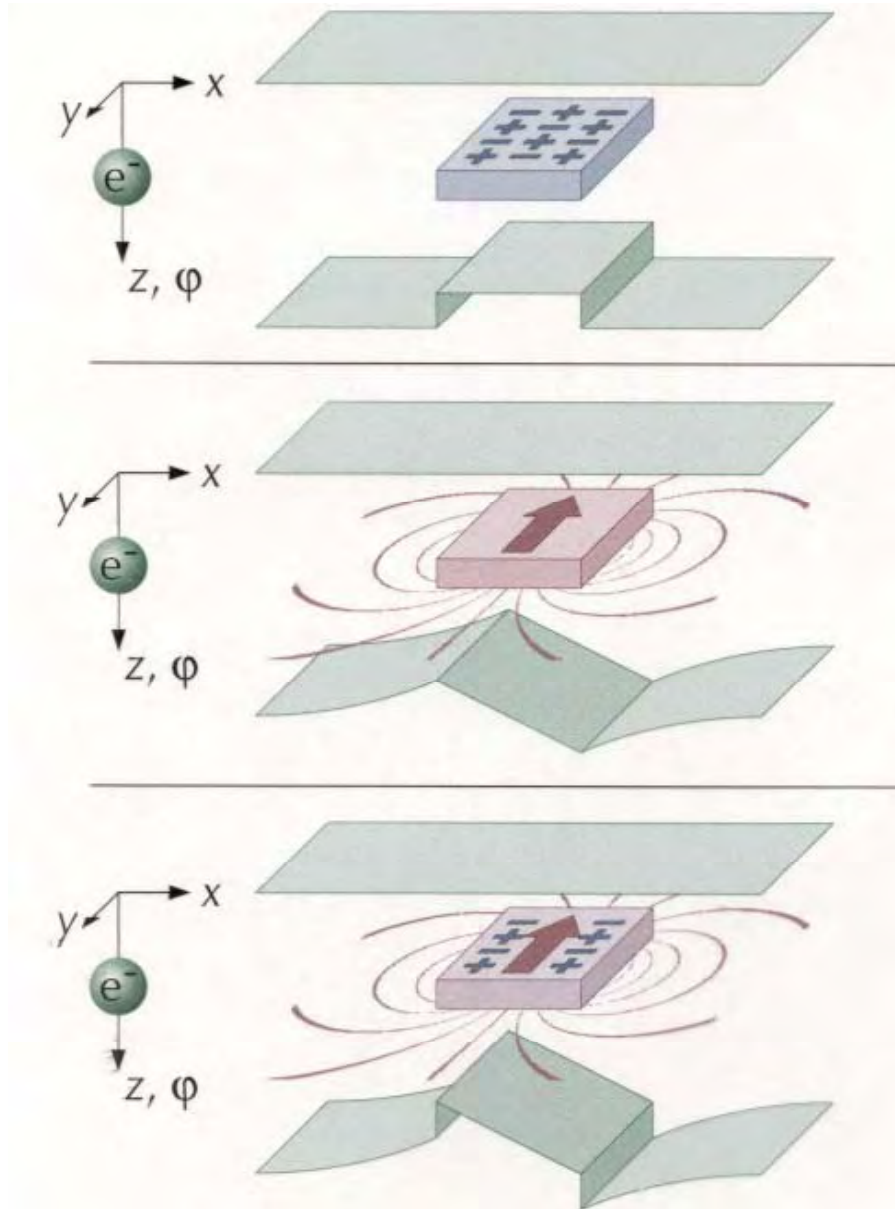
$C_s = 7430 \text{ mm (!!)}$

$C_c = 39 \text{ mm}$

Point resolution $\approx 2 \text{ nm}$

Magnetic samples

Pb. 2: How to separate the magnetic and electrostatic contributions B and $C_E V_i t$?



Phase shift induced by the local electrostatic potential:

$$\varphi_{elect}(\mathbf{x}) = C_E \int V(x, z) dz$$

Phase shift induced by the local magnetic field :

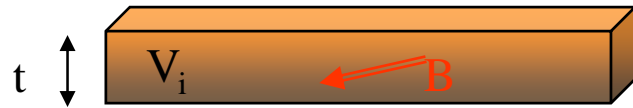
$$\varphi_{mag}(\mathbf{x}) = -\frac{e}{\hbar} \iint B_n(\mathbf{x}, z) dx dz$$

Total phase shift

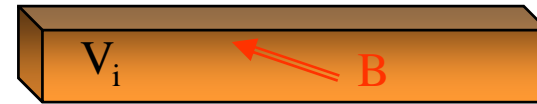
$$\varphi_{Tot}(\mathbf{x}) = \varphi_{elect}(\mathbf{x}) + \varphi_{mag}(\mathbf{x})$$

Pb. 2: How to separate the magnetic and electrostatic contributions B and $C_E V_i t$?

1/ Taking two holograms (4 with the reference holograms) switching upside down the sample
 \Rightarrow The sign of the magnetic contribution is reversed and the electrostatic one remains.



$$\varphi_1 = C_E V_i t - \frac{e}{\hbar} \iint B_n ds$$



$$\varphi_2 = C_E V_i t + \frac{e}{\hbar} \iint B_n ds$$

$$\varphi_2 - \varphi_1 = 2 \frac{e}{\hbar} \iint B_n ds \quad \varphi_2 + \varphi_1 = 2 C_E V_i t$$

Problem : finding back the same area of interest

2/ Taking two holograms with two different high voltages

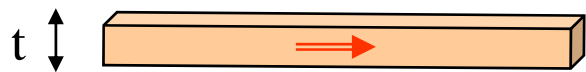
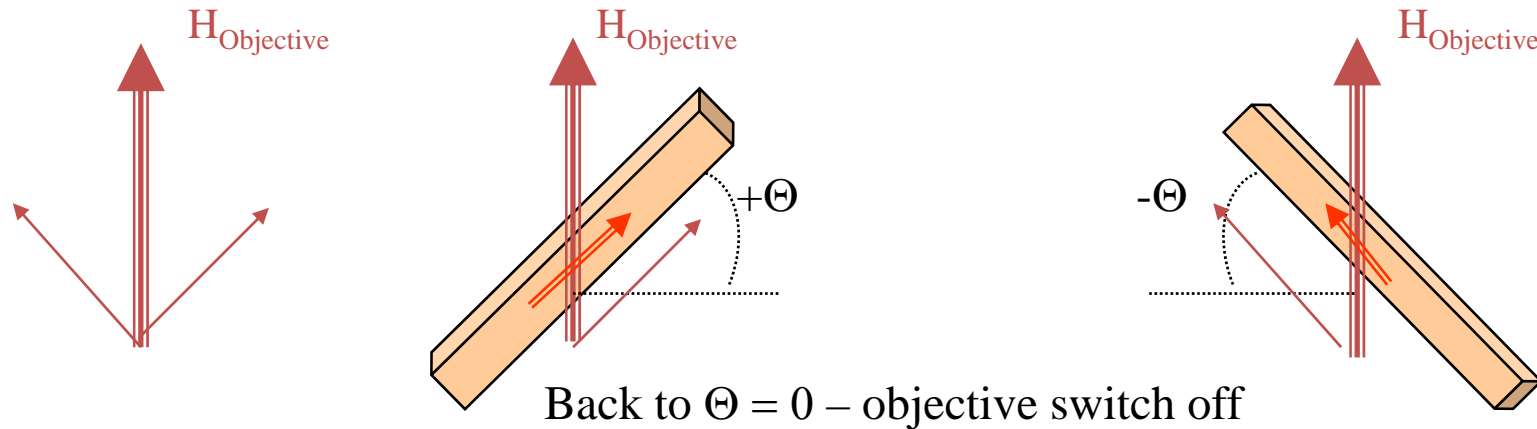
$$C_E = \left(\frac{2\pi}{\lambda} \right) \left(\frac{E + E_0}{E(E + 2E_0)} \right)$$

$$\begin{aligned} \varphi_1 &= C_{E_1} V_i t - \frac{e}{\hbar} \iint B_n ds \\ \varphi_2 &= C_{E_2} V_i t - \frac{e}{\hbar} \iint B_n ds \end{aligned} \quad \left| \right.$$

$$V_i t = \frac{\varphi_2 - \varphi_1}{(C_{E_2} - C_{E_1})}$$

Problem : keeping the same acquisition conditions (magnification, alignment...) when changing the high voltage.

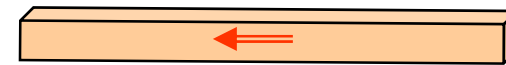
3/ Using the high magnetic field of the objective (~2 Tesla) to saturate the magnetic sample in two opposite directions



$$\varphi_1 = C_E V_i t - \frac{e}{\hbar} \cdot t \cdot \int B_n(x) \cdot dx$$

$$\frac{e}{\hbar} \cdot t \cdot \int B_n(x) \cdot dx = \frac{\varphi_2 - \varphi_1}{2}$$

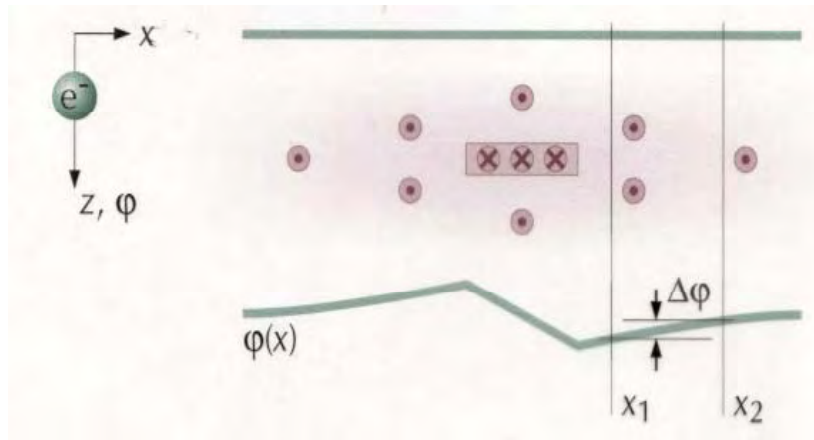
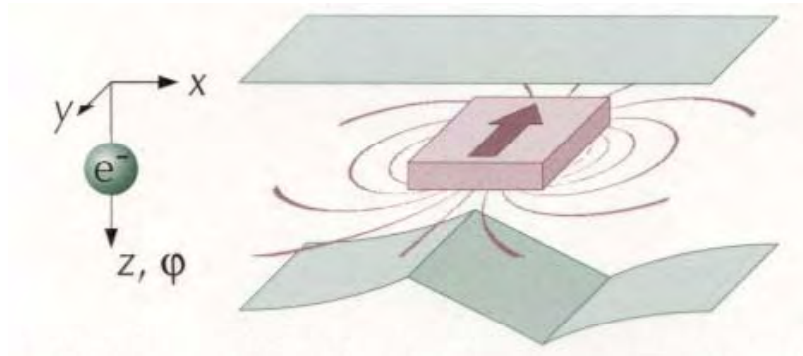
Magnetic contribution to the phase shift



$$\varphi_2 = C_E V_i t + \frac{e}{\hbar} \cdot t \cdot \int B_n(x) \cdot dx$$

$$C_E V_i t = \frac{\varphi_2 + \varphi_1}{2}$$

Electrostatic contribution to the phase shift



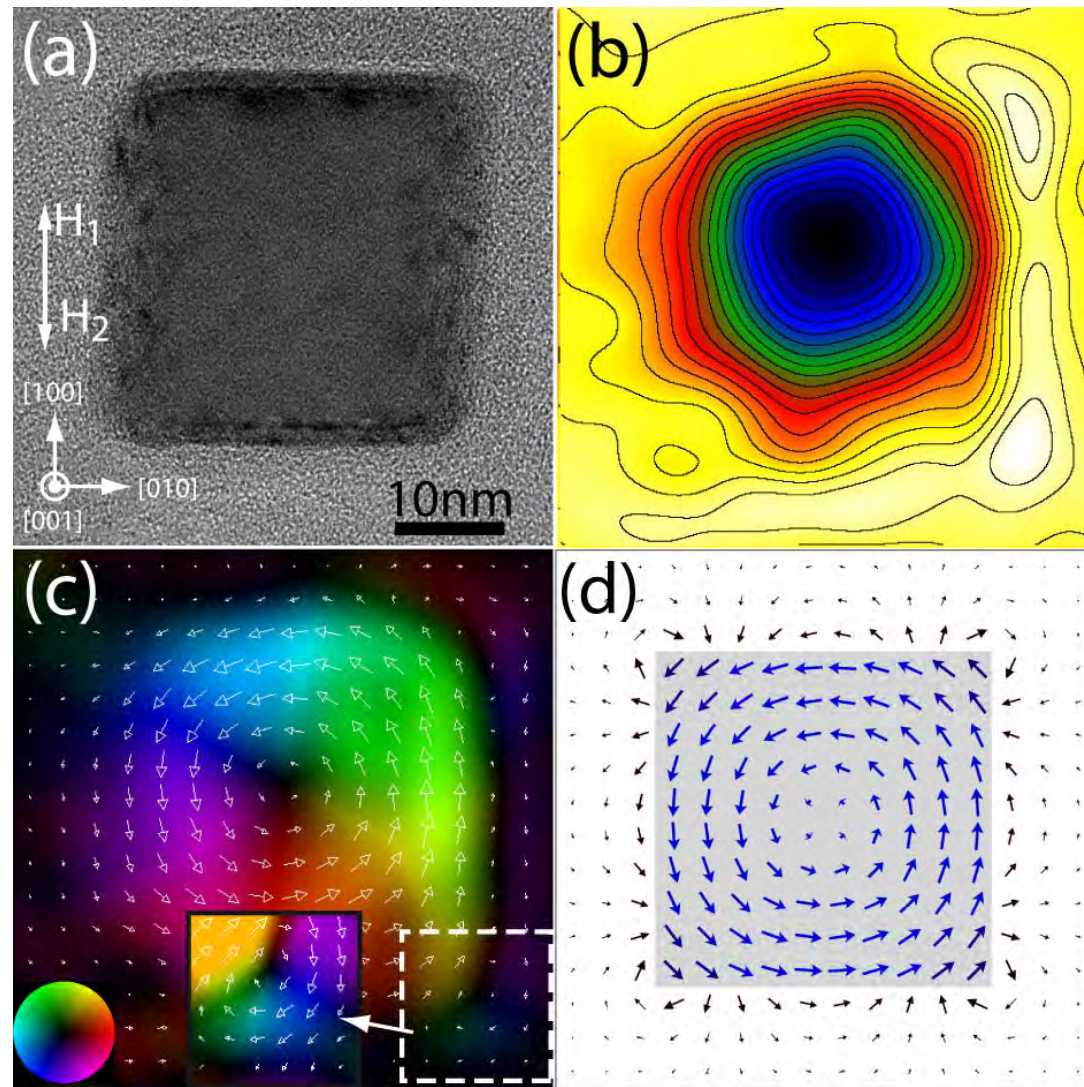
$$\Delta\varphi = \varphi(x_2, y) - \varphi(x_1, y) = -\frac{e}{\hbar} \int_{\xi=x_2}^{x_1} \int B_n(\xi, y, z) d\xi dz$$

$$\frac{\partial\varphi(x)}{\partial x} = \frac{e}{\hbar} B_n(x).t$$

- The gradient of the phase is proportional to the in-plane component of the magnetic induction B_n
- The equiphase contour gives the direction of the magnetic induction B

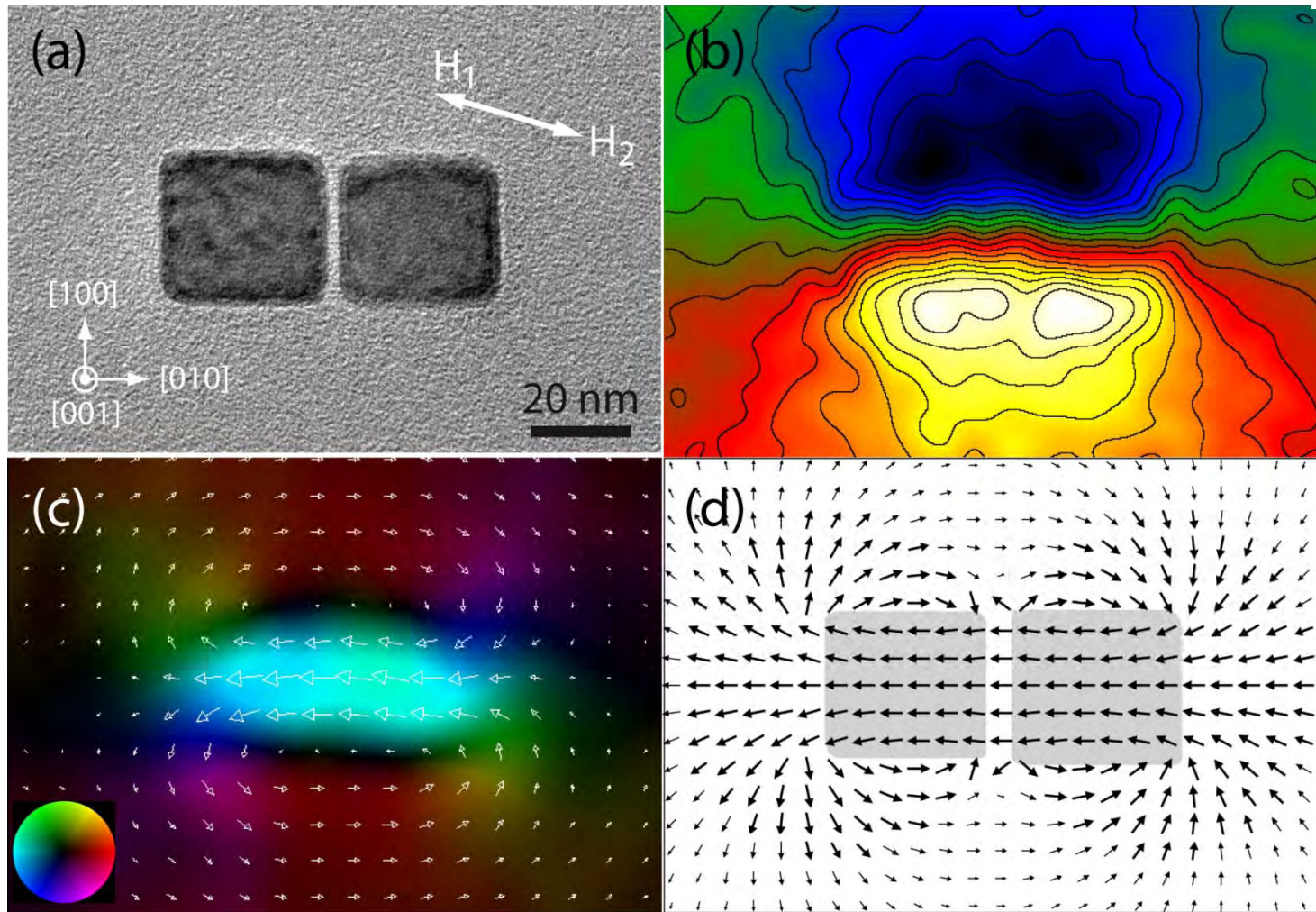
Magnetic Configuration of an isolated Fe Nanocube

Snoeck et al., *Nano Lett.*, 2008, 8 (12), pp 4293–4298



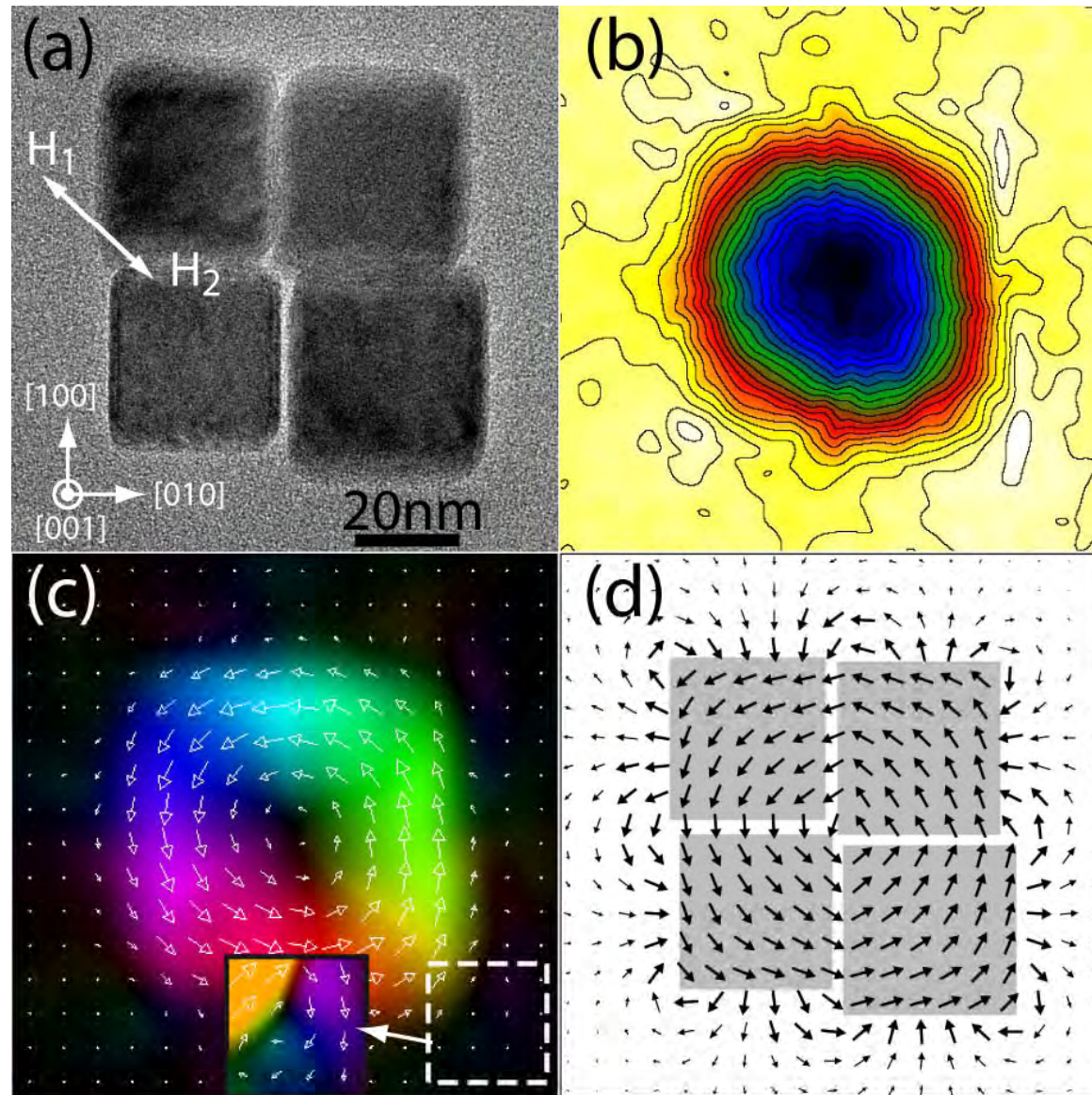
(a) TEM image, (b) Magnetic contribution to the phase shift, (c) Magnetic induction mapping, (d) micromagnetic simulation (OOMMF)

Magnetic Configuration of two neighbouring Fe Nanocubes



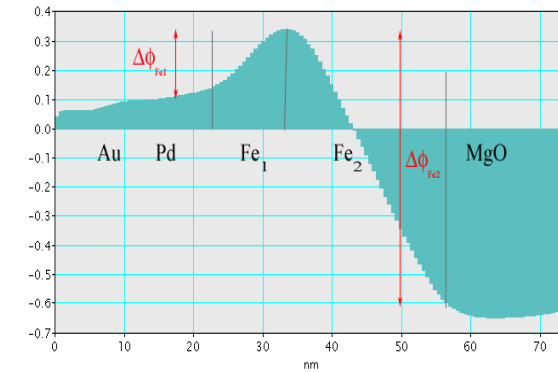
(a) TEM image, (b) Magnetic contribution to the phase shift, (c) Magnetic induction mapping, (d) micromagnetic simulation (OOMMF)

Magnetic Configuration of four neighbouring Fe Nanocubes

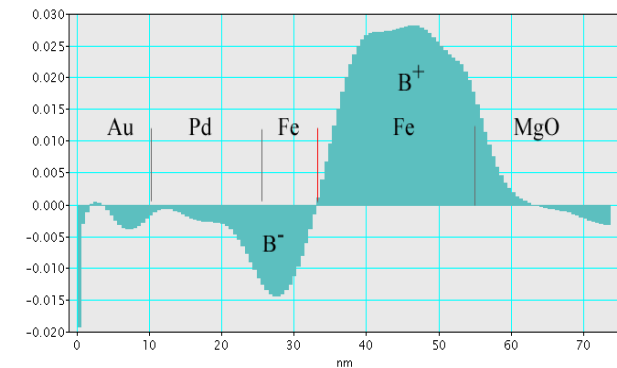
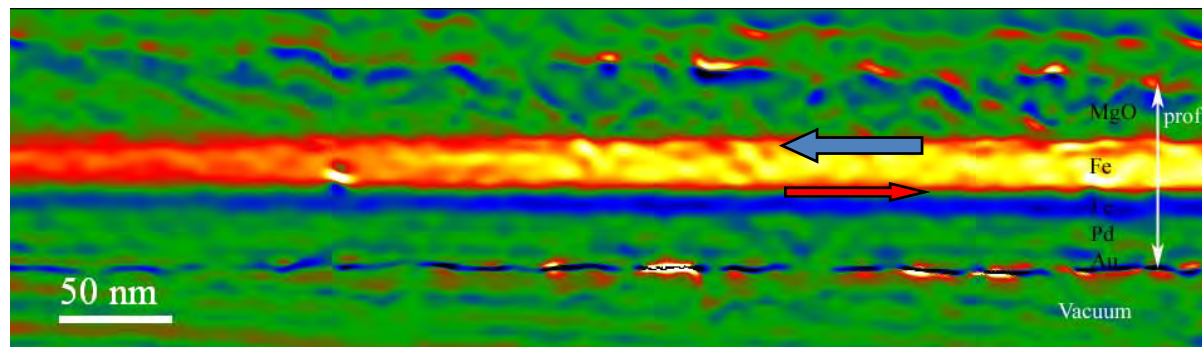


Magnetic multilayers

$$\phi(x) = e/\hbar \int B_{\perp}(x) t(x) dx$$



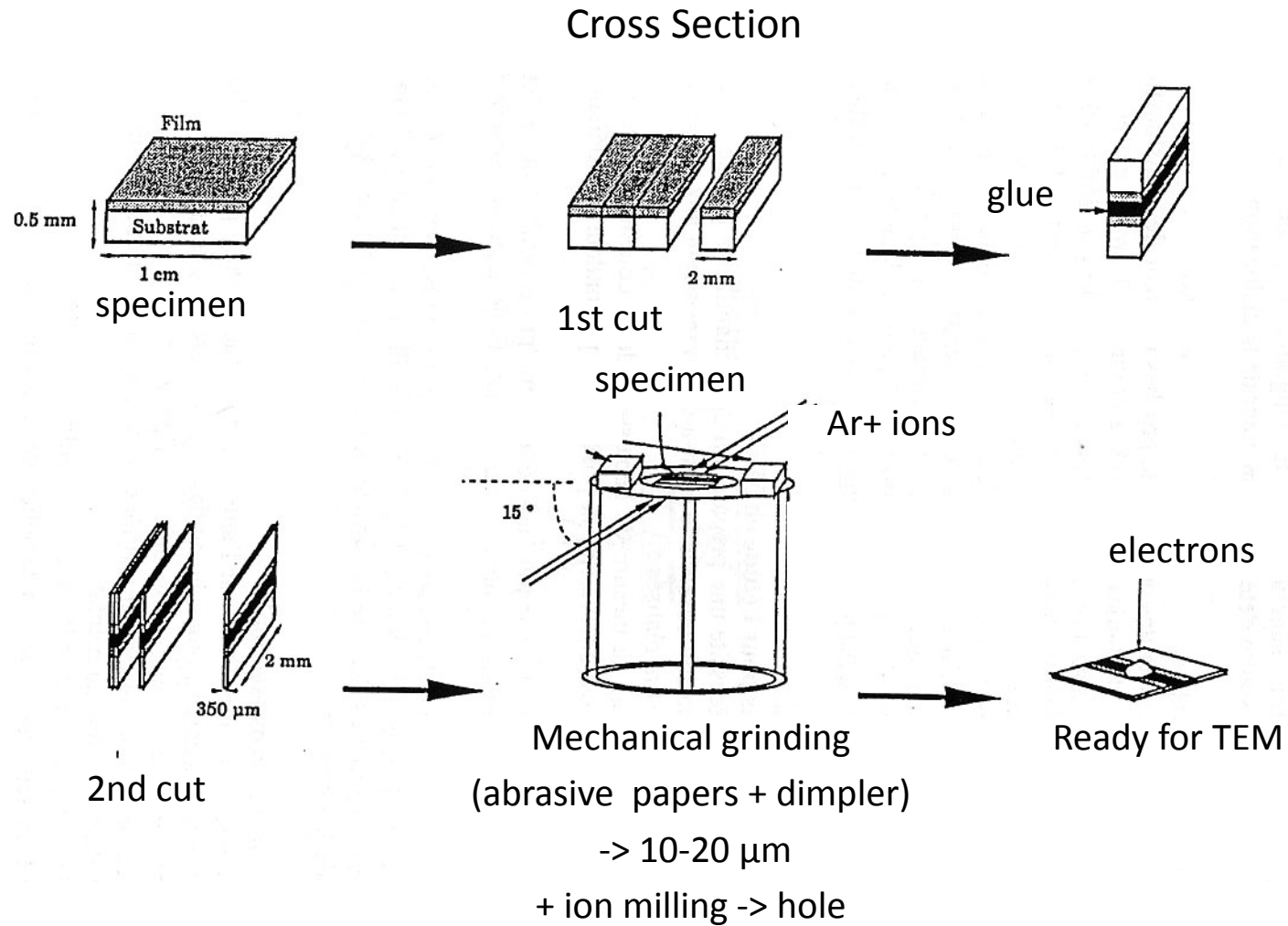
$$d\phi(x)/dx = (e / \hbar) \cdot t \cdot B_{\perp}(x)$$



E. Snoeck, P. Baules, G. BenAssayag, C. Tiusan, F. Greullet, M. Hehn, A. Schuhl

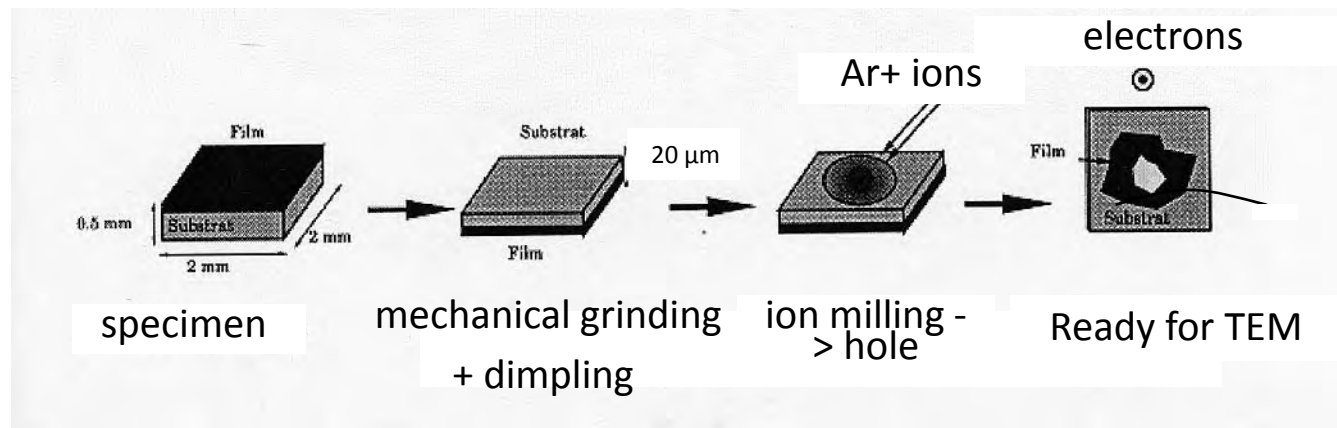
J. Phys.: Condens. Matter 20 (2008) 055219.

Sample preparation

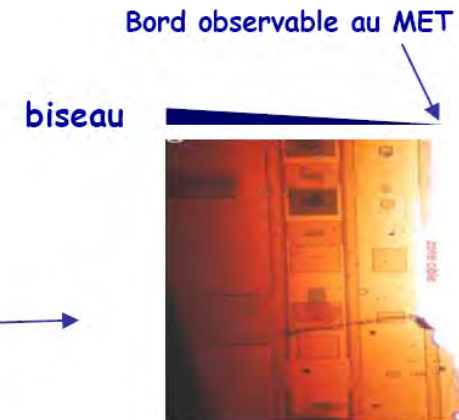
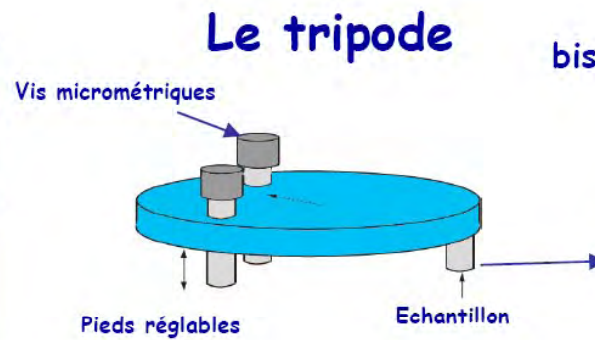
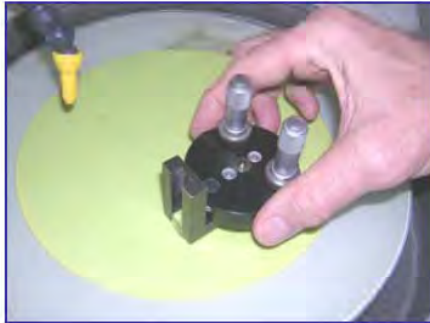


Sample preparation

Plan view



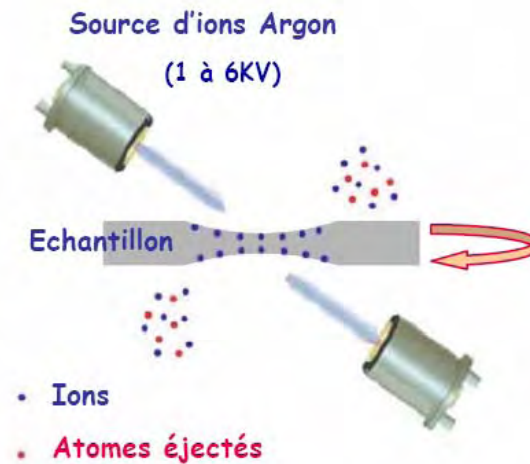
Sample preparation



Amincisseurs ioniques DUAL ION MILL et PIPS (Gatan)

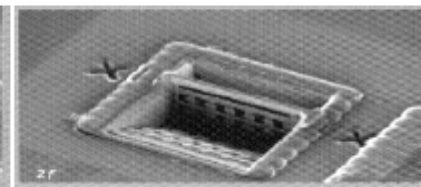
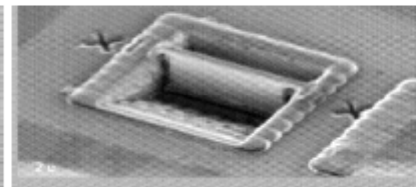
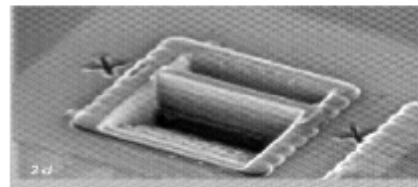
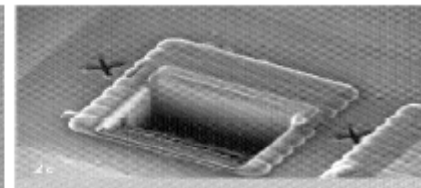
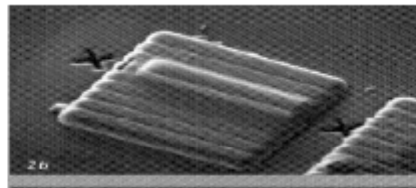
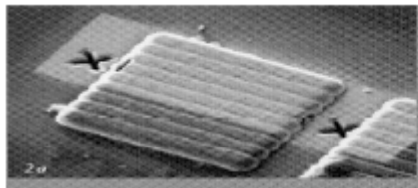
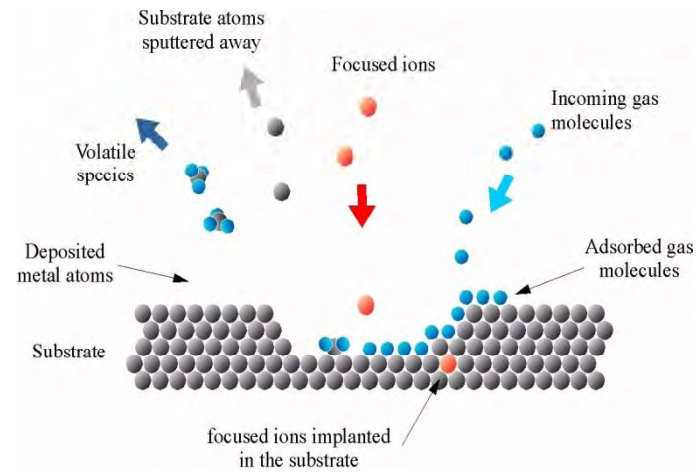


Principe

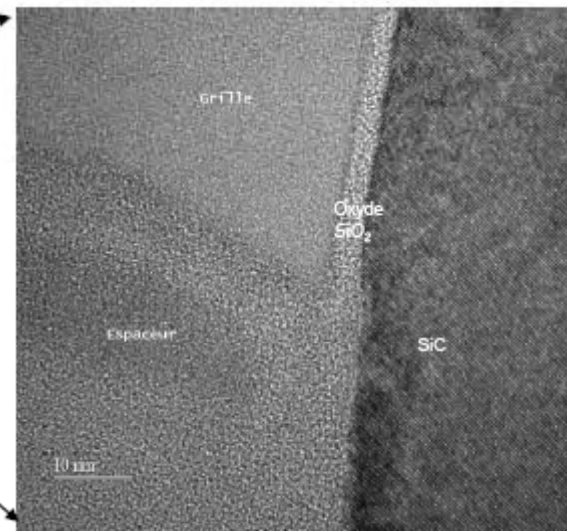
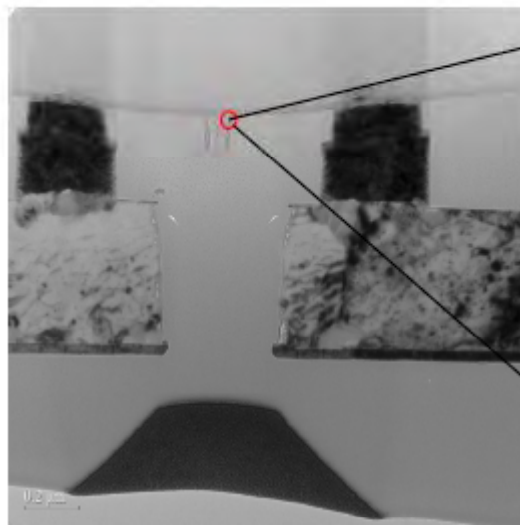
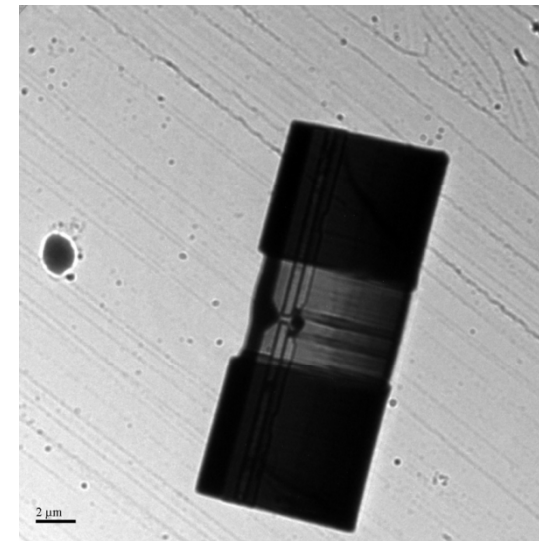
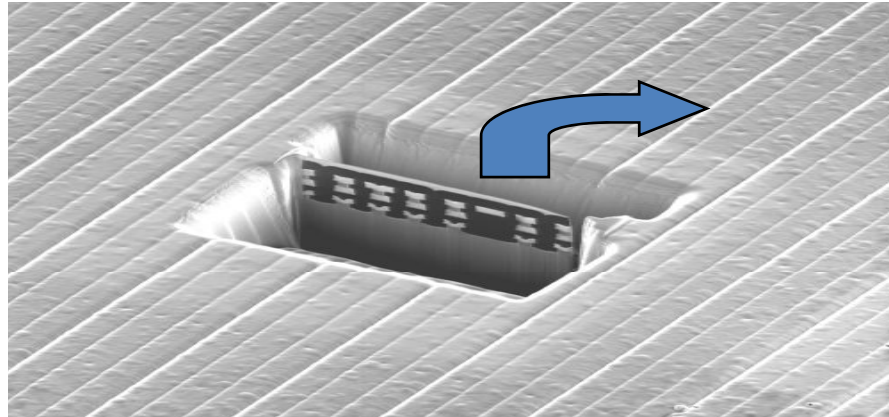


Jacques Crestou
Service préparation

Sample preparation: FIB



Sample preparation: FIB



Summary

Good TEM results from

50% specimen preparation,

20% TEM's price

30% guy's experience...