

Problem Set 3: Solutions

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1.) a.) $V(x)$

(warning: I used q instead of x)

$$i\hbar \dot{S} = \frac{\partial S}{\partial t} \Rightarrow \frac{\partial S}{\partial p} \frac{\partial p}{\partial t} + \frac{\partial S}{\partial q} \frac{\partial q}{\partial t} = i\hbar \dot{S}$$

$$- \frac{\partial S}{\partial p} \frac{\partial H}{\partial q} + \frac{\partial S}{\partial q} \frac{\partial H}{\partial p} = i\hbar \dot{S}$$

$$\{H, S\} = i\hbar \dot{S}$$

Verlet: $e^{i\hbar \dot{S} \Delta t} S(p(0), q(0)) = S(p(\Delta t), q(\Delta t))$

can write: $i\hbar = i\hbar p + i\hbar q$

$$\rightarrow i\hbar p S = \frac{\partial S}{\partial p} \frac{\partial p}{\partial t}$$

$$i\hbar q S = \frac{\partial S}{\partial q} \frac{\partial q}{\partial t}$$

$$e^{i\hbar \dot{S} \Delta t} \sim e^{i\hbar p \frac{\Delta t}{2}} e^{i\hbar q \Delta t} e^{i\hbar p \frac{\Delta t}{2}}$$

$$e^{i\hbar p \frac{\Delta t}{2}} e^{i\hbar q \Delta t} \underbrace{e^{i\hbar p \frac{\Delta t}{2}} S(p(0), q(0))}$$

$$\left[1 + \frac{i\hbar p \Delta t}{2} + \frac{(i\hbar p)^2 \Delta t^2}{2 \cdot 2^2} + \dots \right] S(p(0), q(0))$$

$$i\hbar p S(p(0), q(0)) = \frac{\partial S(p(0), q(0))}{\partial p} \frac{\partial p}{\partial t}$$

$$(i\hbar p)^2 S(p(0), q(0)) = \frac{\partial^2 S(p(0), q(0))}{\partial p^2} \left(\frac{\partial p}{\partial t}\right)^2$$

$$\Rightarrow e^{i\hbar p \frac{\Delta t}{2}} S(p(0), q(0)) = S\left(p(0) + \Delta t \frac{\partial p(0)}{\partial t}, q(0)\right)$$

$$= S\left(p(0) + \frac{\Delta t}{2} \hat{F}(0), q(0)\right) = S\left(p\left(\frac{\Delta t}{2}\right), q(0)\right)$$

Similarly: $e^{i L_a \Delta t} g(p(\frac{\Delta t}{2}), q(0)) = g(p(\frac{\Delta t}{2}), q(\Delta t))$

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Verlet algorithm:

- start with $p(0), q(0)$

- using $q(0) \rightarrow F(q(0))$, calculate $p(\frac{\Delta t}{2})$

$$p(\frac{\Delta t}{2}) = q(0) + \frac{\Delta t}{2} F(q(0))$$

- using $p(\frac{\Delta t}{2}) \rightarrow$ calculate $q(\Delta t)$

$$q(\Delta t) = q(0) + \frac{p(\frac{\Delta t}{2})}{m} \Delta t$$

- using $q(\Delta t) \rightarrow$ calculate $p(\Delta t) = p(\frac{\Delta t}{2}) + \frac{\Delta t}{2} F(q(\Delta t))$

b.) $V(x, x) + \frac{p^2}{2m} + \frac{P^2}{2M} = H(p, P, x, X)$

$$g(p, P, x, X) \rightarrow g(p(0), P(0), x(0), X(0))$$

$$i L_S = \frac{\partial g}{\partial t} \frac{\partial p}{\partial t} + \frac{\partial g}{\partial P} \frac{\partial P}{\partial t} + \frac{\partial g}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial g}{\partial X} \frac{\partial X}{\partial t}$$

$$i L_m = \frac{\partial g}{\partial P} \frac{\partial P}{\partial t} + \frac{\partial g}{\partial X} \frac{\partial X}{\partial t} \rightarrow \text{Lagrangian for small mass particle}$$

$$i L_M = \frac{\partial g}{\partial p} \frac{\partial p}{\partial t} + \frac{\partial g}{\partial x} \frac{\partial x}{\partial t} \rightarrow \text{Lagrangian for large mass particle}$$

$$e^{i L \Delta t} \approx e^{i L_m \frac{\Delta t}{2}} \underbrace{e^{i L_m \Delta t}} e^{i L_M \frac{\Delta t}{2}}$$

$$e^{i L_m \Delta t} = [e^{i L_m \Delta t}]^n$$

$$\cancel{e^{i L_m \Delta t}}$$

Now apply usual breakup:

$$i L_m = i L_m^{(p)} + i L_m^{(x)}$$

$$i L_M = i L_M^{(p)} + i L_M^{(x)}$$

Start with: $\delta(p(0), \rho(0), x(0), X(0))$

$$e^{iL_n \frac{\Delta t}{2}} \sim e^{iL_n^{(K)} \frac{\Delta t}{4}} e^{iL_n^{(P)} \frac{\Delta t}{2}} e^{iL_n^{(X)} \frac{\Delta t}{2}}$$

$$e^{iL_m \Delta t} \sim e^{iL_m^{(K)} \frac{\Delta t}{2}} e^{iL_m^{(P)} \Delta t} e^{iL_m^{(X)} \frac{\Delta t}{2}}$$

propagation:

$$e^{iL_n \frac{\Delta t}{2}} e^{iL_m \Delta t} e^{iL_n \frac{\Delta t}{2}} \delta(p(0), \rho(0), x(0), X(0))$$

$$= e^{iL_n \frac{\Delta t}{2}} \underbrace{e^{iL_m \Delta t}}_{\text{propagated in } n \text{ small steps } (\Delta t)}$$

$$\delta(p(\Delta t), \rho(\frac{\Delta t}{2}), x(\Delta t), X(\frac{\Delta t}{2}))$$

$$\vdots$$

$$\delta(p(\Delta t), \rho(\frac{\Delta t}{2}), x(\Delta t), X(\frac{\Delta t}{2}))$$

apply $e^{iL_n \frac{\Delta t}{2}}$ (propagate heavy degrees of freedom)

$$\rightarrow \delta(p(\Delta t), \rho(\Delta t), x(\Delta t), X(\Delta t))$$

Note: the small step propagations were done such that the small degrees of freedom ~~was~~ propagated but the heavy one was stationary.

Problem 2:

a) $P(x) T(x|x') = P(x') T(x'|x)$ (detailed balance)
symmetric matrix $M(x|x')$

$$T(x|x') = M(x|x') A(x|x')$$

$$\rightarrow P(x) A(x|x') = P(x') A(x'|x)$$

$$A(x|x') = \min \left[1, \frac{P(x')}{P(x)} \right] \quad (\text{acceptance probability})$$

b.) matrix $M(x|x')$ is not symmetric

$$P(x) M(x|x') A(x|x') = P(x') M(x'|x) A(x'|x)$$

$$A(x|x') = \min \left[1, \frac{P(x') M(x'|x)}{P(x) M(x|x')} \right]$$

check! $P(x') M(x'|x) > P(x) M(x|x')$

in this case

$$A(x|x') = 1$$

$$A(x'|x) = \frac{P(x) M(x|x')}{P(x') M(x'|x)}$$

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$$P(x) M(x|x') = P(x') M(x'|x) \frac{P(x) M(x|x')}{P(x') M(x'|x)} \checkmark$$

or $P(x') M(x'|x) < P(x) M(x|x')$
 $\rightarrow A(x'|x) = 1 \quad A(x|x') = \frac{P(x') M(x'|x)}{P(x) M(x|x')}$

$$\rightarrow \cancel{P(x) M(x|x')} = P(x)$$

\rightarrow detailed balance gives $\frac{P(x) M(x|x') P(x') M(x'|x)}{P(x') M(x'|x)} = P(x) M(x|x')$