

MATH544: APPLIED MATHEMATICS II

Lecture 4

- 1) Parabolic Equations (Heat eqn).
- 2) The Maximum principle for the heat eqn and consequences
- 3) Green's function for the Heat eqn.
- 4) Heat flow in an infinite rod.

Parabolic equations

1. Initial and Boundary value problems for a finite interval

$$\left. \begin{array}{l} u_t = a u_{xx}, \quad t > 0 \quad 0 < x < L \quad (a > 0) \\ (BV) \quad u(0, t) = u(L, t) = 0, \quad t \geq 0 \\ (IV) \quad u(x, 0) = f(x), \quad 0 \leq x \leq L \end{array} \right\} \quad (1)$$

Formal solution:

since the PDE and the boundary conditions are homogeneous
then we can use the method of separation of variables

$$u(x, t) = X(x) T(t) \quad (2)$$

Then using this ansatz in the heat equation we get

$$\frac{T'}{T} = a \frac{X''}{X} \quad (3)$$

so that $X(0) = X(L) = 0$. From (3) we get
two equations

$$T' = -\lambda^2 T, \quad X'' + \lambda^2 X = 0 \quad (4)$$

The reason that (3) is equal to $-\lambda^2 a$ depends
on the boundary conditions. You can show that this
is the only choice. The eigenvalue problem

$$X'' + \lambda^2 X = 0, \quad 0 < x < L \quad (5)$$

$$X(0) = X(L) = 0$$

has a nontrivial solution if and only if (3) is equal to $-\alpha\lambda^2$.
The solution of the above problem is ($\lambda = k_n$).

$$X(x) = A_n \sin(k_n x), \quad k_n = \frac{n\pi}{L}, \quad n=1, 2, \dots \quad (6)$$

and the function $T(t)$ is easily found as

$$T(t) = e^{-\alpha k_n^2 t} \quad (7)$$

Hence

$$u_n(x, t) = X_n(x) T_n(t) = A_n e^{-\alpha k_n^2 t} \sin(k_n x), \quad (8)$$

$n=1, 2, \dots$. The heat equation with the boundary conditions has infinitely many solutions $u_n(x, t)$ in (8). To satisfy the initial condition we use the method of superposition, hence

$$u(x, t) = \sum_{n=1}^{\infty} A_n e^{-\alpha k_n^2 t} \sin(k_n x), \quad t > 0 \quad (9)$$

at $t=0$

$$u(x, 0) = \sum_{n=1}^{\infty} A_n \sin(k_n x) = f(x) \quad (10)$$

$$\Rightarrow A_n = \frac{2}{L} \int_0^L f(x) \sin(k_n x) dx \quad (11)$$

(3)

Then $u(x,t)$ in (9) is the formal solution
of the above problem in (1) in $D = \{(x,t) | t > \frac{x}{2}, 0 \leq x \leq L\}$.

- Justification of the problem (of the solution):

Term by term differentiation of the expression in (9)
requires uniform convergence

$$u(x,t) \Rightarrow \sum_{n=1}^{\infty} A_n e^{-ak_n^2 t} \sin(k_n x) \quad (12)$$

$$u_{xx}(x,t) \Rightarrow -a \sum_{n=1}^{\infty} k_n^2 A_n e^{-ak_n^2 t} \sin(k_n x). \quad (13)$$

$$u_t(x,t) \Rightarrow -a \sum_{n=1}^{\infty} k_n^2 A_n e^{-ak_n^2 t} \sin(k_n x) \quad (14)$$

Hence the uniform convergence of the series in the
last expression requires:

$$\left| \sum k_n^2 A_n e^{-ak_n^2 t} \sin(k_n x) \right| \leq \sum k_n^2 |A_n| e^{-ak_n^2 t} \quad (15)$$

It is enough to have $|A_n| \leq M$ where
 M is any positive constant. Then

$$\left| \sum k_n^2 A_n e^{-ak_n^2 t} \sin(k_n x) \right| \leq M \sum_{n=1}^{\infty} k_n^2 e^{-ak_n^2 t} \quad (16)$$

$$< M \sum_{n=1}^{\infty} k_n^2 e^{-ak_n^2 \frac{x}{2}} \quad (17)$$

using the Weierstrass M-test one can show that (4)

$$\sum_{n=1}^{\infty} k_n^2 e^{-ak_n^2 t}, \quad a > 0 \quad (18)$$

Hence.

$$\left| \sum k_n^2 A_n e^{-ak_n^2 t} \sin(k_n x) \right| < \infty. \quad (19)$$

for all $x \in [0, L]$ and $t \geq \tilde{t}/2 > 0$. Hence $u(x, t)$ in (9) with $|A_n| \leq M \quad \forall n = 1, 2, \dots$ is the solution of the problem (1) in $x \in [0, L]$ and $t \geq \tilde{t}/2 > 0$. The condition $|A_n| \leq M$ requires

~~for $t > 0$ we have to show that the series in (10) converges also uniformly.~~

a condition on the initial value $f(x)$. from (11)

$$|A_n| \leq \frac{2}{L} \int_0^L |f(x)| dx \quad (20)$$

Since it is enough to assume that $f(x)$ is continuous in $[0, L]$. Define M so that

$$|f(x)| \leq \frac{M}{2} \quad \forall x \in [0, L]. \quad (21)$$

we get

$$|A_n| \leq M. \quad (22)$$

(5)

For $t=0$ we have to show that the series in (10) converges also uniformly. For this purpose we need $A_n \sim \frac{1}{n^2}$. Hence using the integration by part we get

$$\begin{aligned} A_n &= \frac{2}{L} \int_0^L f(x) \sin(k_n x) dx \\ &= \frac{2}{L} \left[-\frac{1}{k_n} f(x) \cos(k_n x) \right]_0^L + \frac{1}{k_n} \int_0^L f'(x) \cos k_n x dx \\ &= \frac{2}{L k_n} \int_0^L f'(x) \cos k_n x dx \end{aligned} \quad (23)$$

$$f(0) = f(L) = 0 \quad (24)$$

$$\begin{aligned} A_n &= \frac{2}{L k_n} \int_0^L f'(x) \cos k_n x dx \\ &= \frac{2}{L k_n} \left[\frac{1}{k_n} f'(x) \sin k_n x \right]_0^L - \frac{1}{k_n} \int_0^L f''(x) \sin k_n x dx \\ &= -\frac{2}{L k_n^2} \int_0^L f''(x) \sin k_n x dx , \end{aligned} \quad (25)$$

$$\Rightarrow |A_n| \leq \frac{2}{L k_n^2} \int_0^L |f''(x)| dx , \quad (26)$$

(6)

let

$$M_2 = \int_0^L |f''(x)| dx \Rightarrow \quad (27)$$

$$|A_n| \leq \frac{2}{Lk_n^2} M_2 = \frac{2M_2 L}{n^2 \pi^2} = \frac{1}{n^2} \left(\frac{2M_2 L}{\pi^2} \right) \quad (28)$$

Then we have the following theorem.

Theorem: If $f(x)$ is bounded and integrable in $(0, L)$
then the series $\sum_{n=1}^{\infty} A_n e^{-ak_n t} \sin(k_n x)$, $k_n = \frac{n\pi}{L}$, $n=1, 2, \dots$

yields ~~a unique~~ solution of the heat equation.

$u_f = u(x, 0)$ with $u(0, t) = u(L, t) = 0$ for $t > 0$
and $0 \leq x \leq L$. If in addition $f(x)$ has a

continuous second derivative and $f(0) = f(L) = 0$

then $u(x, t) = \sum_{n=1}^{\infty} A_n e^{-ak_n^2 t} \sin k_n x$ is uniformly

convergent for $t \geq 0$ on $0 \leq x \leq L$ and also

satisfies the initial condition $u(x, 0) = f(x)$.

(7)

2. Heat equation with source (inhomogeneous heat eqn).

Suppose now that heat sources or sinks are present in the rod. Then the initial and boundary value problem for the rod becomes

$$u_t = \alpha u_{xx} + F(x, t), \quad t > 0, \quad x \in (0, L) \quad (27)$$

$$\left. \begin{array}{l} u(0, t) = 0, \quad u(L, t) = 0, \quad t \geq 0 \\ u(x, 0) = f(x), \quad 0 \leq x \leq L \end{array} \right\} \quad (28)$$

where $F(x, t)$ is a given function of x and t (source density). We assume that, from our experience of the previous problem,

$$u(x, t) = \sum_{n=1}^{\infty} u_n(t) e^{-\alpha k_n^2 t} \sin(k_n x) \quad (29)$$

where $k_n = \frac{n\pi}{L}$, $n = 1, 2, \dots$. To have a solution the source function $F(x, t)$ must have a similar series expansion

$$F(x, t) = \sum_{n=1}^{\infty} F_n(t) e^{-\alpha k_n^2 t} \sin(k_n x) \quad (30)$$

Then

$$F_n(t) e^{-\alpha k_n^2 t} = \frac{2}{L} \int_0^L F(x, t) \sin(k_n x) dx \quad (31)$$

for all $n = 1, 2, \dots$

Substituting (31) into the heat equation in (29) we get (8)

$$u_n'(t) = F_n(t), \quad n=1, 2, \dots \quad (33).$$

with the solution

$$u_n(t) = u_n(0) + \int_0^t F_n(c) dc. \quad (34)$$

for all $n=1, 2, \dots$ with , from (30)

$$\sum_{n=1}^{\infty} u_n(0) \sin(k_n x) = f(x) \quad (35)$$

or

$$A_n = u_n(0) = \frac{2}{L} \int_0^L f(x) \sin(k_n x) dx, \quad n=1, 2, \dots \quad (36)$$

Hence the formal solution is

$$u(x, t) = \sum_{n=1}^{\infty} A_n e^{-ak_n^2 t} \sin(k_n x) + \sum_{n=1}^{\infty} \left[\int_0^t F_n(c) dc \right] e^{-ak_n^2 t} \sin(k_n x) \quad (37)$$

If $f(x)$ is bounded and integrable in $(0, L)$ then the first term series in (37) is uniformly convergent.

Furthermore, the Fourier series in (35) is uniformly convergent if $f(0)=f(L)=0$ and $f(x)$ has continuous second derivatives in $0 \leq x \leq L$.

The second term in (37): $F(0,t) = F(L,t) = 0$ from (9)
 its property in (31). The Fourier series in (32)
 is uniformly convergent if $F(x,t)$ has continuous partial
 derivative wrt x in $(0,L)$ then

$$F_n(t) e^{-ak_n^2 t} = -\frac{2}{Lk_n^2} \int_0^L F''(x,t) \sin(k_n x) dx, \quad n=1,2,\dots \quad (38)$$

Then

$$|F_n(t) e^{-ak_n^2 t}| \leq \frac{2}{Lk_n^2} \int_0^L |F''(x,t)| dx, \quad n=1,2,\dots \quad (39)$$

since $F''(x,t) = F_{xx}(x,t)$ is continuous in $(0,L)$, $t > 0$
 we have $|F_{xx}(x,t)| < M \quad \forall x \in [0,L] \Rightarrow$

$$|F_n(t) e^{-ak_n^2 t}| \leq \frac{2M}{k_n^2} = \frac{1}{n^2} \left(\frac{2ML^2}{\pi^2} \right) \quad n=1,2,\dots \quad (40)$$

Then

$$\sum_{n=1}^{\infty} F_n(t) e^{-ak_n^2 t} < \infty \quad (41)$$

consequently Weierstrass M-test guarantees the
 uniform convergence of the series in the RHS of (31)

(10)

Term by term differentiation, for instance u_{xx} :

First term in (32) is alright (we have studied it in the first problem). We shall study the second term in (32), $(u_t) = \text{first} + \text{second}$.

$$\text{Second} = \sum \left[F_n(t) e^{-ak_n^2 t} \sin k_n x - \left(\int_0^t F_n(\bar{c}) d\bar{c} \right) a k_n^2 e^{-ak_n^2 t} \sin(k_n x) \right] \quad (42)$$

$$\begin{aligned} |\text{Second}| &\leq \sum \left[|F_n(t)| e^{-ak_n^2 t} + a k_n^2 \int_0^t |F_n(c)| d\bar{c} e^{-ak_n^2 t} \right] \\ &\leq \frac{2ML^2}{\pi^2} \sum \frac{1}{n^2} + \sum a k_n^2 \left(\int_0^t |F_n(c)| e^{-ak_n^2 c} dc \right) e^{ak_n^2 (t-c)} \end{aligned}$$

$$\leq \frac{2ML^2}{6} + \sum a k_n^2 \frac{2M}{k_n^2} \int_0^t e^{ak_n^2 (c-t)} dc.$$

Here we used (40).

$$\begin{aligned} |\text{Second}| &\leq \frac{2ML^2}{6} + 2Ma \sum \int_0^t e^{ak_n^2 (c-t)} dc \\ &\leq \frac{2ML^2}{6} + 2Ma \sum \frac{1}{ak_n^2} (1 - e^{-ak_n^2 t}) \\ &\leq \frac{2ML^2}{6} + \frac{2ML^2}{\pi^2} \sum \frac{1}{n^2} = \frac{2ML^2}{3} < \infty \end{aligned}$$

(11)

Weierstrass M-test implies the uniform convergence of the series (42) and hence RHS of (37). Then we have the following theorem.

Theorem: Suppose $f(x)$ is bounded and integrable in $0 \leq x \leq L$ and $F(x, t)$ has continuous second order partial derivatives with respect to x , with $F(0, t) = F(L, t) = 0, t \geq 0$. Then the series (30) with $u_n(t)$ given in (34) is a solution of the inhomogeneous heat equation and the boundary conditions in (29). If in addition $f(x)$ has continuous second derivative in $(0, L)$ and $f'(0) = f'(L) = 0$ then (30) also satisfies the initial condition in (29).

3) Heat equation with Inhomogeneous boundary conditions

Suppose that the previous initial and boundary value problem becomes

$$u_t = a u_{xx} + f(x, t), \quad 0 < x < L, \quad t > 0 \quad (a > 0)$$

$$\left. \begin{array}{l} u(0, t) = g(t), \quad u(L, t) = h(t), \quad t > 0 \\ u(x, 0) = f(x) \end{array} \right\} (43)$$

where $g(t)$ and $h(t)$ are some differentiable functions of t . We can transform the above problem to the initial value problem with homogeneous boundary conditions by letting

$$v(x, t) = u(x, t) + \frac{L-x}{L} g(t) + \frac{x}{L} h(t) \quad (44)$$

which leads to

$$\left. \begin{array}{l} v_t = a v_{xx} + \bar{f}(x, t) \quad 0 < x < L, \quad t > 0 \\ v(0, t) = v(L, t) = 0, \quad t > 0 \\ v(x, 0) = \tilde{f}(x) \quad 0 \leq x \leq L \end{array} \right\} (45)$$

where

$$\left. \begin{array}{l} \bar{f}(x, t) = f(x, t) - \frac{L-x}{L} g(t) - \frac{x}{L} h(t) \\ \tilde{f}(x) = f(x) - \frac{L-x}{L} g(0) - \frac{x}{L} h(0) \end{array} \right\} (46)$$

Hence $v(x, t)$ satisfies our previous problem.

Discussions

a) It is clear from the solution of the homogeneous heat equation

$$u(x,t) = \sum A_n e^{-ak_n^2 t} \sin(k_n x)$$

where

$$A_n = \frac{2}{L} \int_0^L f(x) \sin k_n x$$

with

$$|A_n| \leq 2 \max_{x \in [0,L]} |f(x)|$$

and

$$|u(x,t)| \leq \max |f(x)| \gamma$$

where $\gamma = 2 \sum e^{-ak_n^2 t}$. Indeed we shall show later that $|u(x,t)| \leq \max_{x \in [0,0.1]} |f(x)|$.

b) Solution is unique: This is clear from the above expression. Given data f we have a unique solution

(14)

c) continuously dependence on the data (stability of the solution)

$$|u_1 - u_2| \leq \sum |A_n^{(1)} - A_n^{(2)}| e^{-\alpha k_n^2 t} \leq |A_n^{(1)} - A_n^{(2)}| e^{-\alpha k_n^2 \tilde{t}/2}$$

$$\leq \sum |A_n^{(1)} - A_n^{(2)}| , \quad t \geq \tilde{t}/2 > 0$$

but

$$A_n^{(1)} - A_n^{(2)} = \frac{2}{L k_n^2} \int_0^L (f_1'' - f_2'') \sin k_n x dx$$

$$|A_n^{(1)} - A_n^{(2)}| \leq \frac{2}{L k_n^2} \max |f_1'' - f_2''|$$

$$\Rightarrow |u_1 - u_2| \leq \frac{2L^2}{\pi^2} \max |f_1'' - f_2''| \frac{\pi^2}{6}$$

$$\leq \frac{L^2}{3} \max |f_1'' - f_2''|$$

This proves the stability of the initial
and boundary value problem

d) Internal energy

$$E(t) = \rho c \int_0^L u(x,t) dx$$

$$\text{as } t \rightarrow \infty \quad u(x,t) \rightarrow 0 \rightarrow E(t) \rightarrow 0$$

$$\frac{dE}{dt} = \rho c \int_0^L \frac{\partial u}{\partial t} dx = \rho c \int_0^L u_{xx} dx$$

$$= \rho c \left[u_x \right]_0^L = \rho c (u_x(L,t) - u_x(0,t))$$

$$u_x = \sum A_n e^{-ak_n^2 t} k_n \cos k_n x$$

$$\begin{aligned} u_x(L,t) - u_x(0,t) &= \sum A_n e^{-ak_n^2 t} k_n (-1)^n - 1 \\ &= -2 \sum_{n=\text{odd}} A_n k_n e^{-ak_n^2 t} \end{aligned}$$

$$\frac{dE}{dt} \leq 0 \quad t \geq T/2 > 0$$

e) If $f(x) \geq 0$ the initial temperature then
 $u(x,t) \geq 0$ for all $t \geq 0, 0 \leq x \leq L$

$$\text{and } \lim_{t \rightarrow \infty} u(x,t) = 0$$

Let $a > 0$

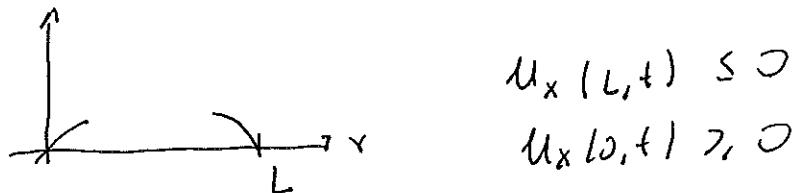
$$u_t - a u_{xx} = 0 \quad , \quad 0 < x < L, \quad t > 0$$

$$u(0, t) = 0, \quad u(L, t) = 0, \quad t > 0$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq L$$

i) $0 \leq u(x, t) \leq \max_{x \in [0, L]} f(x)$

ii) since $u(0, t) = u(L, t) = 0$ then



iii) $E(t) = \int_0^L u(x, t) dx \geq 0, \quad E(0) = 0$
if $u(x, 0) = 0$

$$\frac{dE}{dt} = a \int_0^L u_{xx}(x, t) dx = a[u_x(L, t) - u_x(0, t)] \leq 0$$

$$= -a\gamma(t), \quad \gamma(t) = -u_x(0, t) + u_x(L, t) \geq 0$$

$$E(t) = -a \int_0^t \gamma(u) dt \leq 0 \quad \text{contradiction}$$

$$E(0) = \int_0^L f(x) dx \quad \text{where } u(x, 0) = f(x)$$

$$\Rightarrow E(t) = E(0) - a \int_0^t \gamma(u) dt > 0.$$

(16)

f) If $F \neq 0$ (a non-zero source).

Assume

$$\lim_{t \rightarrow \infty} u(x, t) = u(x)$$

Then what is the limit

$$\lim_{t \rightarrow \infty} u(x, t) \stackrel{?}{=} u(x)$$

stationary state $u_t = 0$

$$a u_{xx} + F = 0 \quad u(0) = u(L) = 0$$

$$u(x, t) = \sum_{n=1}^{\infty} A_n e^{-ak_n^2 t} \sin k_n x + \sum_{n=1}^{\infty} \left[\int_0^t F_n(c) dc \right] e^{-ak_n^2 t} \sin k_n x$$

where

$$F_n(t) e^{-ak_n^2 t} = \frac{2}{L} \int_0^L F(x, t) \sin k_n x dx$$

$$\lim_{t \rightarrow \infty} u(x, t) = \sum \lim_{t \rightarrow \infty} \left(\int_0^t F_n(c) dc / e^{-ak_n^2 t} \right) \sin k_n x$$

$$= \sum \frac{F_n(0)}{ak_n^2 e^{ak_n^2 t}} \sin k_n x$$

(17)

$$\begin{aligned}
 \lim_{t \rightarrow \infty} u(x,t) &= \sum \frac{1}{ak_n^2} F_n(t) e^{-ak_n^2 t} \sin k_n x \\
 &= \sum \frac{1}{ak_n^2} \left(\int_0^L F(x',t) \sin k_n x' dx' \right) \sin k_n x \\
 &= \sum \frac{1}{ak_n^2 L} \lim_{t \rightarrow \infty} \left(\int_0^L F(x',t) \sin k_n x' dx' \right) \sin k_n x \\
 &= \sum \frac{2}{ak_n^2 L} \int_0^L \left(\lim_{t \rightarrow \infty} F(x',t) \right) \sin k_n x' dx' \sin k_n x \\
 &= \sum_{n=1} \frac{2}{ak_n^2 L} \left(\int_0^L F(x') \sin k_n x' dx' \right) \sin k_n x \\
 \text{and } u_{xx} + F &= 0 \\
 u(x,t) &= - \sum_{n=1} \frac{2}{ak_n^2 L} a \int_0^L u_{xx'} \sin k_n x' dx' \sin k_n x \\
 &= - \sum_{n=1} \frac{2}{k_n^2 L} \left[u_x' \Big|_0^L - k_n \int_0^L u_x' \cos k_n x' dx' \right] \sin k_n x \\
 &= \sum \frac{2}{k_n L} \left(\int_0^L u_x' \cos k_n x' dx' \right) \sin k_n x \\
 &= \sum \frac{2}{k_n L} \left[-u \Big|_0^L + k_n \int_0^L u \sin k_n x' dx' \right] \sin k_n x \\
 &= + \sum \frac{2}{L} \left(\int_0^L u \sin k_n x' dx' \right) \sin k_n x \\
 &= \sum_{n=1} \theta_n \sin k_n x = u(x) \quad \checkmark
 \end{aligned}$$

The Maximum Principle of the heat equation and its consequences.

Let

$$u_t - a u_{xx} = F(x, t) \leq 0$$

$$u(0, t) = g(t) \leq 0$$

$$u(l, t) = h(t) \leq 0$$

$$u(x, 0) = f(x) \leq 0$$

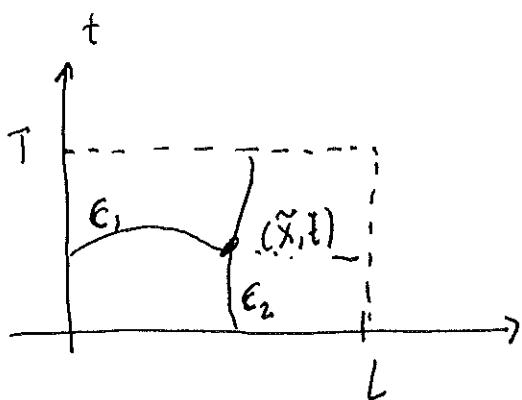
we expect that $u(x, t) \leq 0 \quad \forall 0 \leq x \leq L, t > 0$

We shall prove this: First we prove a proposition

proposition: suppose that $w(x, t)$ satisfies the differential inequality $w_t - a w_{xx} < 0$ in $0 < x < L$,

$0 < t \leq T$, where $a > 0$. Then $w(x, t)$ can not assume a local maximum value at any point in the rectangle $0 < x < L, 0 < t \leq T$

proof. Assume the contrary: let $w(\tilde{x}, \tilde{t})$ be a local maximum value for w at (\tilde{x}, \tilde{t}) ,
 $(0 < \tilde{x} < L, 0 < \tilde{t} \leq T)$



a) Along the curve C_1 :

$w(x, \tilde{t})$, ($0 < x < L$) has a local maximum at $x = \tilde{x}$, hence
 $w_{xx} < 0$ at (\tilde{x}, \tilde{t})

b) Along the curve C_2 : $w(\tilde{x}, t)$, ($0 < t \leq T$) has a local maximum value at $t = \tilde{t}$, hence
 $w_t \geq 0$ at (\tilde{x}, \tilde{t}) (strict inequality occurs at $\tilde{t} = T$)

From (a) and (b) we have $w_t - \alpha w_{xx} \geq 0$ at (\tilde{x}, \tilde{t}) which is a contradiction. Hence there exist no point in the rectangle $0 \leq x \leq L$, $0 \leq t \leq T$, so that $w(x, t)$ assume a local maximum value.

Theorem: Let $\alpha > 0$ and suppose that $u(x, t)$ is continuous in $0 \leq x \leq L$, $0 \leq t \leq T$ (rectangle with its boundaries, $R \cup \partial R = \bar{R}$) and satisfies

$$u_t - \alpha u_{xx} \leq 0 \quad 0 < x < L, \quad 0 < t \leq T$$

$$u(0, t) \leq 0, \quad u(L, t) \leq 0, \quad 0 \leq t \leq T$$

$$u(x, 0) \leq 0, \quad 0 \leq x \leq L$$

Then $u(x, t) \leq 0$ for $(x, t) \in R \cup \partial R$

(20)

Proof: By contradiction: Assume to the contrary that u is positive at some point, say $u(\tilde{x}, \tilde{t}) > 0$

Define

$$w(x, t) = u(x, t) - \varepsilon t$$

where $\varepsilon > 0$ and $0 \leq x \leq L$, $0 \leq t \leq T$. Then for any choice of ε we find:

$$w_t - a w_{xx} = u_t - a u_{xx} - \varepsilon < 0$$

$$w(0, t) = u(0, t) - \varepsilon t \leq 0 \quad 0 \leq t \leq T.$$

$$w(L, t) = u(L, t) - \varepsilon t \leq 0$$

$$w(x, 0) = u(x, 0) \leq 0 \quad 0 \leq x \leq L$$

previous lemma and the above boundary conditions show that for every $\varepsilon > 0$ $w(x, t)$ can not have a (positive) maximum in the closed interval $0 \leq x \leq L$, $0 \leq t \leq T$. On the other hand, for small enough $\varepsilon > 0$ we have

$$w(\tilde{x}, \tilde{t}) = u(\tilde{x}, \tilde{t}) - \varepsilon \tilde{t} > 0$$

but this is a contradiction. Hence there exists no point (\tilde{x}, \tilde{t}) in R so that $u(\tilde{x}, \tilde{t}) > 0$

(21)

Corollary 3. let $a > 0$ and suppose that $u(x,t)$ is continuous in $0 \leq x \leq L$, $0 \leq t \leq T$ and satisfies

$$u_t - a u_{xx} \geq 0 \quad \text{for } 0 < x < L, \quad 0 < t \leq T$$

$$u(0,t) \geq 0, \quad u(L,t) \geq 0 \quad \text{in } 0 \leq t \leq T$$

$$u(x,0) \geq 0 \quad \text{in } 0 \leq x \leq L$$

$$\text{Then } u(x,t) \geq 0 \quad \text{for } 0 \leq x \leq L, \quad 0 \leq t \leq T$$

proof: Take $\sqrt{-u}$. Use theorem the previous thm. for $\sqrt{-u}$.

Corollary 2 let $a > 0$, suppose $u(x,t)$ is defined and continuous in $0 \leq x \leq L$, $0 \leq t \leq T$ and satisfies

$$u_t = a u_{xx} \quad 0 < x < L, \quad 0 < t \leq T$$

$$\text{let } M = \max \left\{ \max_{x \in [0,L]} u(x,0), \max_{0 \leq t \leq T} u(0,t), \max_{0 \leq t \leq T} u(L,t) \right\}$$

$$m = \min \left\{ \min_{x \in [0,L]} u(x,0), \min_{0 \leq t \leq T} u(0,t), \min_{0 \leq t \leq T} u(L,t) \right\}$$

$$\text{Then } m \leq u(x,t) \leq M \quad \text{for all } 0 \leq x \leq L, \quad 0 \leq t \leq T.$$

Proof. i) let $v(x,t) = u(x,t) - M$. Then:

$$v_t - v_{xx} = 0 \quad 0 < x < L, \quad 0 < t \leq T$$

$$v(0,t) \leq 0, \quad v(L,t) \leq 0 \quad 0 < t \leq T.$$

$$v(x,0) \leq 0 \quad 0 \leq x \leq L$$

$$\Rightarrow v(x,t) \leq 0 \Rightarrow u(x,t) \leq M.$$

ii) let $w(x,t) = M - u(x,t)$. Then

$$w_t - w_{xx} = 0$$

$$w(0,t) \leq 0, \quad w(L,t) \leq 0 \quad 0 \leq t \leq T$$

$$w(x,0) \leq 0 \quad 0 \leq x \leq L$$

$$\Rightarrow w(x,t) \leq 0 \Rightarrow u(x,t) \geq M.$$

Hence we obtain $M \leq u(x,t) \leq M$

for all $0 \leq x \leq L, \quad 0 \leq t \leq T$.

Corollary 3. Let $a > 0$, suppose that $u(x, t)$ defined and continuous in $0 \leq x \leq L$, $0 \leq t \leq T$ and satisfies

$$u_t - a u_{xx} = F(x, t) \quad 0 < x < L, \quad 0 < t < T$$

and $F(x, t)$ is bounded by a constant N that

$$|F(x, t)| \leq N \quad \text{for all } 0 \leq x \leq L, \quad 0 \leq t \leq T.$$

Let

$$M = \max \left\{ \max_{x \in [0, L]} |u(x, 0)|, \max_{0 \leq t \leq T} |u(0, t)|, \max_{0 \leq t \leq T} |u(L, t)| \right\}$$

Then:

$$|u(x, t)| \leq M + tN \quad \text{for all } 0 \leq x \leq L, \quad 0 \leq t \leq T$$

Proof: i) Let $v(x, t) = u(x, t) - (M + tN)$

$$v_t - a v_{xx} = u_t - a u_{xx} - N = F - N \leq 0$$

$$v(0, t) \leq 0, \quad v(L, t) \leq 0 \quad 0 \leq t \leq T$$

$$v(x, 0) \leq 0 \quad 0 \leq x \leq L$$

$$\Rightarrow v(x, t) \leq 0 \Rightarrow u(x, t) \leq M + tN$$

ii) Let $w(x, t) = -u(x, t) - (M + tN)$

(24)

Then

$$w_t - a w_{xx} = -F - N \leq 0$$

$$w(0, t) \leq 0, \quad w(L, t) \leq 0 \quad 0 \leq t \leq T$$

$$w(x, 0) \leq 0$$

$$\Rightarrow w(x, t) \leq 0$$

$$\Rightarrow -M - tN \leq u(x, t)$$

As a result

$$|u(x, t)| \leq M + tN \leq M + TN$$

Theorem: There exist at most one function $u(x, t)$
 which is continuous in $0 \leq x \leq L, 0 \leq t \leq T$

and satisfies

$$u_t - a u_{xx} = F(x, t) \quad 0 < x < L, \quad t > 0$$

$$u(0, t) = g(t), \quad u(L, t) = h(t), \quad t \geq 0$$

$$u(x, 0) = f(x) \quad 0 \leq x \leq L$$

where f, g, h are given continuous functions (data)

proof: Let $u_1 - u_2 = w$. Then we have

u_1 and u_2 are the solutions of the above
 initial and boundary value problem with the
 same data. Hence it satisfies

(25)

$$w_t - \alpha w_{xx} = 0 \quad 0 < x < l, \quad t > 0$$

$$w(0, t) = w(l, t) = 0 \quad t > 0$$

$$w(x, 0) = 0 \quad 0 \leq x \leq l$$

This means that $m = M = 0$ (corollary).
 $(m \leq w \leq M)$ Hence $w(x, t) = 0$ which leads to $u_1 = u_2$.

Theorem. Let $\alpha > 0$, suppose that $u(x, t)$ is continuous on $0 \leq x \leq l$, $0 \leq t \leq T$ and satisfies the initial and boundary value problem.

$$\frac{\partial}{\partial t} u - \alpha \frac{\partial^2}{\partial x^2} u = F \quad 0 < x < l, \quad 0 < t \leq T$$

$$u(0, t) = g(t), \quad u(l, t) = h(t), \quad 0 \leq t \leq T$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq l$$

Then this initial and boundary value problem is stable, i.e., the solutions continuously depend on the data.

proof: let the data $(F_1, f_1, g_1, h_1) \rightarrow u_1(x, t)$ and $(F_2, f_2, g_2, h_2) \rightarrow u_2(x, t)$, i.e

(26)

$$\frac{\partial}{\partial t} u_i - \alpha \frac{\partial^2}{\partial x^2} u_i = F_i(x, t), \quad 0 < x < L, \quad 0 < t \leq T$$

$$u_i(0, t) = g_i(t), \quad u_i(L, t) = h_i(t), \quad 0 \leq t \leq T$$

$$u_i(x, 0) = f_i(x), \quad 0 \leq x \leq L$$

for $i=1, 2$. Suppose further that

$$|f_1 - f_2| < \varepsilon, \quad |g_1 - g_2| < \varepsilon$$

$$|h_1 - h_2| < \varepsilon, \quad |F_1 - F_2| < \varepsilon$$

for all $0 \leq x \leq t, \quad 0 \leq t \leq T$. Then

using the Lemma Corollary 3 we get

$$|u_1 - u_2| \leq (1+T)\varepsilon. \quad (N=M=\varepsilon)$$

Another proof of the uniqueness theorem in the previous theorem.

proof: let $w(x,t) = u_1 - u_2$ where u_1 and u_2

are two different solutions of the initial and boundary value problem. Then in the same data, then

$$w_t - a w_{xx} = 0 \quad 0 < x < L, t > 0$$

$$w(0,t) = w(L,t) = 0 \quad t > 0$$

$$w(x,0) = 0, \quad 0 \leq x \leq L$$

Define now the energy functional

$$E(t) = \int_0^L w^2 dx \geq 0$$

with $E(0) = 0$, and

$$\begin{aligned} \frac{dE}{dt} &= \int_0^L 2w w_t dx = 2a \int_0^L w w_{xx} dx \\ &= 2a \left[w w_x \right]_0^L - \int_0^L w_x^2 dx = -2a \int_0^L w_x^2 dx \leq 0 \end{aligned}$$

hence $E' \leq 0$ for all $t > 0$. Hence we have

$$E(t) \geq 0 \quad \forall t > 0$$

$$E(0) = 0$$

$$E'(t) \leq 0 \quad \forall t > 0$$

(24'')

Hence the only possibility is that $E(t) = 0$ $\forall t \geq 0$. This means that

$$w(x, t) = 0 \quad \forall x \in (0, L), t \geq 0$$

$$\text{or } u_1(x, t) = u_2(x, t) \quad \forall x \in [0, L], t \geq 0.$$

This proves the uniqueness.

The Green's Function (of the heat equation)

We have studied that the initial value problem

$$u_t = a u_{xx} \quad 0 \leq x \leq L, t > 0$$

$$u(x, 0) = f(x) \quad 0 \leq x \leq L \quad (1)$$

$$u(0, t) = u(L, t) = 0 \quad t > 0$$

has the solution

$$u(x, t) = \sum_{n=1}^{\infty} c_n e^{-a\lambda_n^2 t} \sin(\lambda_n x), \quad \lambda_n = \frac{n\pi}{L} \quad (1)$$

where

$$c_n = \frac{2}{L} \int_0^L f(y) \sin(\lambda_n y) dy \quad (2)$$

we assumed that $f(0) = f(L) = 0$ and $f(x)$ is twice differentiable function. For $t > 0$ we insert (2) in (1)

$$u(x, t) = \sum_{n=1}^{\infty} \frac{2}{L} \int_0^L f(y) \sin(\lambda_n y) dy e^{-a\lambda_n^2 t} \sin(\lambda_n x)$$

$$= \int_0^L \left[\frac{2}{L} \sum_{n=1}^{\infty} e^{-a\lambda_n^2 t} \sin \lambda_n y \sin \lambda_n x \right] f(y) dy$$

\uparrow converges uniformly and all derivatives are uniformly bounded on $0 \leq x \leq L, t > t_0$ for any $t_0 > 0$
hence interchanging sum and integral is legal

Define

$$G(x, y, t) = \frac{2}{L} \sum_{n=1}^{\infty} e^{-\alpha n^2 t} \sin \lambda_n y \sin \lambda_n x$$

for $0 \leq x, y \leq L$ and $t > 0$

$$\Rightarrow u(x, t) = \int_0^L G(x, y, t) f(y) dy.$$

$$\text{at } t=0 \quad G(x, y, 0) = \delta(x-y).$$

$G(x, y, t)$ is the GF for the heat equation.

$$\text{i) } G_t(x, y, t) = \alpha G_{xx}(x, y, t) = \alpha G_{yy}(x, y, t)$$

for $0 \leq x, y \leq L, t > 0$

$$\text{ii) } G(0, y, t) = G(L, y, t) = 0$$

for $0 \leq y \leq L, t > 0$

$$\text{iii) } G(x, y, t) = G(y, x, t)$$

for $0 \leq x, y \leq L, t > 0$

$$\text{iv) } u_{(x, t)} \rightarrow (x_0, t) \left(\int_0^{t_0} \{ G(x, y, t) \}_{y=t} dy \right) = f(x_0)$$

$$0 \leq x_0 \leq L$$

$$u(x,t) = \int_0^L G(x,y,t) f(y) dy$$

$u(x,t)$ satisfies term by term differentiation of $G(x,y,t)$ the heat equation and satisfies the boundary condit. $f(x)$ has a continuous second derivative

sequence of

Given twice differentiable functions $\{f_n(x)\}$ with $f_n(0) = f_n(L) = 0$. Then given $\varepsilon > 0$ there is a number $N(\varepsilon)$ such that

$$n, m > N(\varepsilon) \Rightarrow |f_n(x) - f_m(x)| < \varepsilon \text{ for } 0 \leq x \leq L$$

let $u_n(x,t)$ be the solution to the heat equation with zero boundary condit. and initial temperature $f_n(x)$. By the maximum principle

$$n, m > N(\varepsilon) \text{ implies that } |u_n(x,t) - u_m(x,t)| < \varepsilon$$

for $0 \leq x \leq L$, $t > 0$, and consequently the sequence $\{u_n(x,t)\}$ converges uniformly to a function $u(x,t)$ on $0 \leq x \leq L$, $t > 0$.

Evidently $u(x,t)$ is continuous. Moreover for each (x,t) with $t > 0$

$$\begin{aligned} u(x,t) &= \lim_{n \rightarrow \infty} u_n(x,t) = \lim_{n \rightarrow \infty} \int_0^L G(x,y,t) f_n(y) dy \\ &= \int_0^L G(x,y,t) f(y) dy \end{aligned}$$

Since $u(x,t)$ is continuous on $0 \leq x \leq L$, $t > 0$ and \exists :

$$u(x_0, 0) = f(x_0)$$

Then we have

Theorem: Let the initial temperature $f(x)$ be continuous on $0 \leq x \leq L$ and satisfy $f(0) = f(L) = 0$. Then initial and boundary value problem (Q) has a unique solution, continuous solution defined on $0 \leq x \leq L$, $t \geq 0$. This solution is given by

$$u(x,t) = \int_0^L G(x,y,t) f(y) dy$$

for $t > 0$

Heat flow in infinite and semi-infinite Rods

The temperature $u(x,t)$ in an infinite rod, taken to be the entire x -axis is given by

$$\begin{aligned} u_t &= a u_{xx} & -\infty < x < \infty, t > 0 \\ u(x,0) &= f(x) & -\infty < x < \infty \end{aligned}$$

where $f(x)$ is the initial temperature distribution.
We use Fourier transform

$$\hat{g}(\omega) = \int_{-\infty}^{\infty} e^{i\omega x} g(x) dx$$

Inverse transform

$$g(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-i\omega x} \hat{g}(\omega) d\omega$$

(6F5)

multiply the heat equa~~l~~^s by $e^{i\omega x}$ and (31)

$$\int_{-\infty}^{\infty} (u_t - a u_{xx}) e^{i\omega x} dx = \hat{u}_t - a \int_{-\infty}^{\infty} u_{xx} e^{i\omega x} dx$$

$$= \hat{u}_t + i a \omega \int_{-\infty}^{\infty} u_x e^{i\omega x} dx$$

$$u_x \Big|_{x \rightarrow \infty} = \hat{u}_t + a \omega^2 \hat{u} = 0$$

$$u \Big|_{x \rightarrow \infty} = \hat{u}(\omega, t) = A(\omega) e^{-a\omega^2 t}$$

$$= \int_{-\infty}^{\omega} e^{i\omega x} u(x, t) dx$$

$$\hat{u}(\omega, t) = A(\omega) = \int_{-\infty}^{\omega} e^{i\omega x} f(x) dx$$

 \Rightarrow

$$u(x, t) = \frac{1}{2\pi} \int_0^{\omega} e^{-i\omega x} \hat{u}(\omega, t) d\omega.$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\omega} A(\omega) e^{-a\omega^2 t - i\omega x} d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\omega} \int_{-\infty}^{\omega} e^{i\omega y} f(y) dy e^{-a\omega^2 t - i\omega x} d\omega$$

~~$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\omega} dy f(y) G(x, y, t)$$~~

(6F.6)

(32)

$$u(x,t) = \int_{-\infty}^{\infty} k(x-y, t) f(y) dy$$

where $k(x-y, t) =$

$$u(x,t) = \int_{-\infty}^{\infty} dy f(y) \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-aw^2 - i w(x-y)} dw.$$

use $\int_{-\infty}^{\infty} e^{-iw\beta - w^2/4t} dw = \sqrt{4\pi t} e^{-\alpha\beta^2}$, $\alpha > 0$

with $\alpha = 1/4at$, $\beta = x-y$ for $t > 0$

$$u(x,t) = \frac{1}{\sqrt{4\pi at}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4at} f(y) dy$$

$$= \int_{-\infty}^{\infty} k(x-y, t) f(y) dy$$

with

$$k(x-y, t) = \frac{1}{\sqrt{4\pi t}} e^{-(x-y)^2/4t}$$

for $-\infty < x < \infty$, $t > 0$

function $k(x-y, t)$ is called the "fundamental solution" to the heat equation. It plays the same role for the infinite rod that G.F. does for finite length rod.

We found the formal solution: By change variable we get

$$u(x,t) = \frac{1}{\sqrt{\pi t}} \int_{-\infty}^{\infty} e^{-\xi^2} f(x - \sqrt{4\pi t}\xi) d\xi.$$

Change variable $\xi = \frac{x-y}{\sqrt{4\pi t}}$

$$\text{The solution } u(x,t) = \int_{-\infty}^{\infty} k(x-y,0t) f(y) dy$$

Cannot give a solution to the IVP for all continuous function $f(y)$. There should be restriction on f . This integral diverges if $f(x) = e^{xy}$. This suggests that f should be bounded on the real line.

It is easy to show that

$$k_x(x,t) = -\frac{2x}{4t} k(x,t)$$

$$k_t = k_{xx} = \frac{x^2 - 2t}{4t^2} k(x,t)$$

We have the following lemma

Lemma: Let $0 < t_0 < T$, $R > 0$ and $\gamma > 1$ be given. Then there is a constant N depending on t_0 , T , R and γ such that

$$0 \leq k(x-y,0t) \leq N e^{-\frac{|y|^2}{4Rt}}$$

$$|k_x(x-y, t)| \leq N e^{-\frac{y^2}{4at}}$$

$$|k_x(x-y, t)| = |k_{xx}(x-y, t)| \leq N e^{-\frac{y^2}{4at}}$$

for all $|x| \leq R$, $t_0 \leq t \leq T$

Proof:

$$k(x-y, at) = \frac{1}{\sqrt{4\pi at}} e^{-\frac{(x-y)^2}{4at}} \leq \frac{1}{\sqrt{4\pi at_0}} e^{-\frac{(x-y)^2}{4at_0}}$$

$$\leq \frac{1}{\sqrt{4\pi at_0}} e^{-\frac{-x^2 + 2xy - y^2}{4at}} \leq \frac{1}{\sqrt{4\pi at_0}} e^{\frac{2R|y| - y^2}{4at}} e^{-\frac{y^2}{4at}}$$

$$\leq \frac{1}{\sqrt{4\pi at_0}} e^{\frac{2R|y|}{4at}} e^{-\frac{y^2(1 - \frac{1}{R})}{4at}} e^{-\frac{y^2}{4at}}$$

~~~~~

bounded and continuous

has a max. value =  $N$

$$\leq N e^{-\frac{y^2}{4at}}$$

All other functions  $k_x, k_y, k_{xy}$  are bounded  
the bound  $N$  is the largest for all these cases

Since  $\int_{-\infty}^{\infty} e^{-\frac{y^2}{4at}} dy = \sqrt{4\pi at} < \infty$

hence for any bounded, integrable function  $f(y)$   
the integral

$$\int_{-\infty}^{\infty} k(x-y, at) f(y) dy$$

Can be differentiated under the integral sign once wrt  $t$  and twice wrt  $x$  for  $x \in \mathbb{R}$ ,  $0 < t_0 \leq t \leq T$ . Let  $t_0 \rightarrow 0$  and  $R, T \rightarrow \infty$  to see that here derivative can be taken under the integral sign for  $-\infty < x < \infty$  and  $t > 0$

and since  $h_t(x,t) = h_{xx}(x,t) \Rightarrow$  the soln satisfies the heat eqn. hence we have

**Theorem:** Suppose that the initial temperature distribution is continuous and bounded on  $-\infty < x < \infty$ . Then the diffusion problem has a solution which is continuous on  $-\infty < x < \infty, t > 0$ . For  $t > 0$  the soln. is

$$u(x,t) = \int_{-\infty}^{\infty} k(x-y,at) f(y) dy$$

Indeed we have

$$u(x,t) = \begin{cases} \int_{-\infty}^{\infty} k(x-y,at) f(y) dy & \text{for } x < \infty, t > 0 \\ \text{we showed this part} & \end{cases}$$

$$f(x) \quad -\infty < x < \infty \quad t=0$$

We have to show that the continuity of this soln. we proved the continuity for all  $t > 0$  for  $t=0$  we must show

$$\lim_{(x,t) \rightarrow (x_0,0)} \int_{-\infty}^{\infty} k(x-y,at) f(y) dy = f(x_0)$$

(G10) (36)

Proof:

$$\text{i)} \int_{-\infty}^{\infty} k(x-y, at) dy = 1$$

like GF

$$k(x-y, at) = \frac{1}{\sqrt{4\pi at}} e^{-\frac{(x-y)^2}{4at}}$$

$$\frac{1}{\sqrt{4\pi at}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4at}} dy$$

$$\bar{y} = \frac{x-y}{\sqrt{4at}}$$

$$= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-s^2} ds = 1.$$

$$\text{ii)} \int_{-\infty}^{\infty} k(x-y, at) f(y) dy - f(x)$$

$$= \int_{-\infty}^{\infty} k(x-y, at) [f(y) - f(x)] dy$$

$$= \int_{|y-x|<\delta} k(x-y, at) [f(y) - f(x)] dy \quad S_1$$

$$+ \int_{|y-x|>\delta} k(x-y, at) [f(y) - f(x)] dy \quad S_2$$

Given  $\epsilon > 0$  there  $\delta > 0$  so that  $|f(y) - f(x)| < \epsilon$   
 for  $|y-x| < \delta \Rightarrow$

(64)  
(37)

$$\Rightarrow |S_1| < \varepsilon$$

In  $|y-x_0| > \delta$  we have  $f(y)$  is bounded

$$|f(y)| < M$$

$$\begin{aligned} |S_2| &< \int_{|y-x_0| \geq \delta} h(x-y, t) |f(y) - f(x_0)| dy \\ &< 2M \int_{-\infty}^{x_0 - \delta} h(x-y, t) dy + 2M \int_{x_0 + \delta}^{\infty} h(x-y, t) dy \\ &\leq \frac{2M}{\sqrt{\pi t}} \left[ \int_{-\infty}^{(x_0 - \delta - x)/\sqrt{4at}} e^{-s^2} ds + \int_{(x_0 + \delta - x)/\sqrt{4at}}^{\infty} e^{-s^2} ds \right] \end{aligned}$$

$$\text{with } F = \frac{y-x}{\sqrt{4at}}$$

For  $|x-x_0| < \delta/2$  we have

$$|S_2| \leq \frac{2M}{\sqrt{\pi t}} \left[ \int_{-\infty}^{(x_0 - \delta/2 - x)/\sqrt{4at}} e^{-s^2} ds + \int_{\delta/2/\sqrt{4at}}^{\infty} e^{-s^2} ds \right] < \varepsilon$$

provided that  $t > 0$  is sufficiently small to complete the proof.

$$\Rightarrow \lim_{t \rightarrow 0} \int_{-\infty}^{\infty} h(x-y, t) f(y) dy = f(x)$$

Remark: Let  $\delta_n(x)$  be a mass point  
Assume initially that the entire mass  
of a substance is localized at the point  $x_0$   
and the total mass is  $Q$ . Let  $\delta_n(x)$  be  
a sequence of continuous functions which  
vanish for  $|x-x_0| > \frac{1}{n}$  and satisfy

$$\begin{cases} \delta_n \text{ mass } \delta_n \Rightarrow \\ x_0 \in [x_0 - \frac{1}{n}, x_0] \end{cases}$$

$$\int_{-\infty}^{\infty} \delta_n(x) dx = 1$$

Then  $Q \delta_n(x)$  is the concentration of a  
mass distributed ~~near~~ located at  $x_0$  and with  
total mass  $Q$

Let  $u_n(x, t)$  be the solution with the initial  
data  $Q \delta_n(x)$ , then for  $t > 0$

$$\begin{aligned} u_n(x, t) &= \int_{-\infty}^{\infty} k(x-y, at) Q \delta_n(y) dy \\ &= Q \int_{x_0 - \frac{1}{n}}^{x_0 + \frac{1}{n}} k(x-y, at) \delta_n(y) dy \end{aligned}$$

According to the MVT for integrals

$$\begin{aligned} u_n(x, t) &= k(x-y_n, at) \int_{x_0 - \frac{1}{n}}^{x_0 + \frac{1}{n}} Q \delta_n(y) dy \\ &= Q k(x-y_n, at) \end{aligned}$$

$$\lim_{n \rightarrow \infty} u_n(x, t) = Q k(x-x_0, at)$$

$\Rightarrow k(x-x_0, at)$  represents the concentration of unit mass  
concentrated at  $x_0$  when  $t=0$

Semi-infinite rod:

GF12

(39)

$$u_+(x, t) = u_{xx} = 0 \quad x > 0, t > 0$$

$$u(x, 0) = f(x) \quad x > 0$$

$$u(0, t) = 0 \quad t > 0$$

define

$$u^+(x, t) = \int_{-\infty}^{\infty} h(x-y, at) f(y) dy$$

$$u^-(x, t) = \int_{-\infty}^0 h(x-y, at) f(-y) dy$$

$$= \int_0^{\infty} h(x+y, at) f(y) dy$$

$$u(x, t) = u^+ - u^- = \int_0^{\infty} g(x, y, at) f(y) dy$$

$$g(x, y, at) = h(x-y, at) - h(x+y, at)$$

GF for semi-infinite interval  
 $0 < x < \infty$

ASSIGNED EXERCISES OF MATH544: PDE set 6 April, 2001

### THE HEAT EQUATION

Ronald B Guenther and John W Lee, Partial Differential Equations of Mathematical Physics and Integral Equations

1. Solve the following initial value problem

$$u_t = au_{xx}, \quad (a > 0), t > 0, \quad x \in [0, L], \quad (1)$$

$$u(0, t) = u(L, t) = 0, \quad t \geq 0, \quad (2)$$

$$u(x, 0) = f(x), \quad x \in [0, L] \quad (3)$$

**step 1.** Formal Solution: Using the separation of variables we first let  $u(x, t) = T(t)X(x)$ . Then find that

$$T' = akT, \quad X'' = kX$$

where  $k$  is any constant. It must be a negative number  $k = -\lambda^2$ . we find that

$$T = e^{-a\lambda_n^2 t}, X = \sin(\lambda_n x), n = 1, 2, \dots$$

where  $\lambda_n = \frac{n\pi}{L}$ . Hence

$$u(x, t) = \sum_{n=1}^{\infty} \alpha_n e^{-a\lambda_n^2 t} \sin(\lambda_n x) \quad (4)$$

Here  $\alpha_n$  are constants , will be determined by the initial condition (assuming the  $t$  goes to zero exists which requires the uniform convergence of the sine series)

$$f(x) = \sum_{n=1}^{\infty} \alpha_n \sin(\lambda_n x), \quad (5)$$

with

$$\alpha_n = \frac{2}{L} \int_0^L f(x) \sin(\lambda_n x) dx, \quad n = 1, 2, \dots \quad (6)$$

**step 2.** Is the formal solution a solution of the initial and boundary value problem we started with?. For this the given function must satisfy certain conditions. Term-by term differentiation of the series representing  $u(x, t)$  in (4) can be justified at each point  $(\tilde{x}, \tilde{t})$  with  $\tilde{t} > 0$  and  $0 \leq \tilde{x} \leq L$  by showing that the series in (4) and the series representing  $u_t$ ,  $u_x$  and  $u_{xx}$  that result from term-by-term differentiation , converge uniformly in a region  $D$  which contains the point  $(\tilde{x}, \tilde{t})$ . Such a region is

$$D = \{(x, t) | 0 \leq x \leq L, t \geq \frac{\tilde{t}}{2}\}.$$

We first assume that the function  $f(x)$  is bounded in  $[0, L]$

$$|f(x)| \leq M, \quad \text{for all } x \in [0, L]$$

Then from (6) we have

$$|\alpha_n| \leq 2M$$

and series (4) satisfies

$$\left| \sum_{n=1}^{\infty} \alpha_n e^{-a\lambda_n^2 t} \sin(\lambda_n x) \right| \leq 2M \sum_{n=1}^{\infty} \lambda_n^2 e^{-a\lambda_n^2 \tilde{t}/2}.$$

The series at the right hand side converges uniformly in  $D$ . Similar reasoning applies also for the series representing  $u_t$ ,  $u_x$ , and  $u_{xx}$ . Consequently (4) satisfies the equations (1) and (2).

To guarantee that (4) will also satisfy the initial condition (5), we need further assumptions on  $f(x)$ . For the uniform convergence of the series in (5) we need  $\sum_{n=1}^{\infty} |\alpha_n| < \infty$ . If  $f(0) = f(L) = 0$  then

$$\alpha_n = -\frac{2}{\lambda_n^2 L} \int_0^L f'' \sin(\lambda_n x) dx.$$

Hence , if  $f''$  is continuous on  $0 \leq x \leq$ , then it is bounded say by  $M''$ , and

$$|\alpha_n| \leq \frac{2M''L^2}{\pi^2 n^2}$$

This estimate implies that the series (4) for  $u(x, t)$  is uniformly convergent for  $t \geq 0$  and  $0 \leq x \leq L$ . Since  $u(x, t)$  is the uniform limit of continuous functions, it is continuous. Hence we have the following:

**Theorem.** *If  $f(x)$  is bounded and integrable, the series in (4) yields a solution to (1) and (2) for  $t > 0$  and  $0 \leq x \leq L$ . If in addition  $f$  has continuous second derivative and  $f(0) = f(L) = 0$ , then  $u(x, t)$  defined in (4) is uniformly convergent and continuous for  $t \geq 0$  and  $0 \leq x \leq L$  and also satisfies the initial condition (3).*

2. Solve the following initial and boundary value problem:

$$u_t - au_{xx} = F(x, t), \quad t > 0, \quad 0 < x < L, \quad (7)$$

$$u(0, t) = u(L, t) = 0, \quad t \geq 0, \quad (8)$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq L \quad (9)$$

Complete the method of solution given in class. The method of solution is similar to previous problem.

3. Solve the problem:

$$u_t - au_{xx} = F(x, t), \quad 0 < x < L, \quad t > 0, \quad (10)$$

$$u_x(0, t) = u_x(L, t) = 0, \quad t \geq 0, \quad (11)$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq L. \quad (12)$$

Explain also how you can prove the uniqueness of the solution of this problem.

4. Let  $u(x, t)$  be a solution to the problem

$$u_t - au_{xx} = 0, \quad t > 0, \quad 0 \leq x \leq L, \quad (13)$$

$$u(0, t) = u(L, t) = 0, \quad t \geq 0 \quad (14)$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq L \quad (15)$$

let  $\tilde{u}$  be a solution of the above problem with  $f$  replaced by  $\tilde{f}$ . Denote the respective Fourier coefficients by  $\alpha_n$  nad  $\tilde{\alpha}_n$ . Show  $|\alpha_n - \tilde{\alpha}_n| \leq 2||f'' - \tilde{f}''||/\lambda_n^2$

and hence

$$\|u - \tilde{u}\| \leq 2\|f'' - \tilde{f}''\| \left( \sum_{n=1}^{\infty} \frac{1}{\lambda_n^2} \right)$$

where  $\|h\| = \max|h(x)|$  for  $0 \leq x \leq L$ . Deduce continuous dependence and hence the well-posedness of the problem.

5. For  $T > 0$  let

$$Q_T = \{(x, t) | 0 < x < L, 0 < t \leq T\}$$

(this is exactly the region  $D$  defined in the lecture), be a rectangle with its bottom and two vertical sides removed. Let

$$\bar{Q}_T = \{(x, t) | 0 \leq x \leq L, 0 \leq t \leq T\}.$$

Then  $B_T = \bar{Q}_T - Q_T$  consists of the three missing sides of the original rectangle  $Q_T$ . Suppose that  $w(x, t)$  satisfies the differential inequality

$$w_t - a(x, t)w_{xx} - b(x, t)w_x - c(x, t)w < 0 \quad (x, t) \in Q_T$$

where  $a(x, t) \geq 0$  and  $c(x, t) \leq 0$  in  $Q_T$ . Then  $w(x, t)$  can not achieve a positive local maximum in  $Q_T$

6. Let  $a, b$  and  $c$  be as in the previous problem. Suppose that  $u(x, t)$  is continuous on  $\bar{Q}_T$  and satisfies

$$u_t - a(x, t)u_{xx} - b(x, t)u_x - c(x, t)u \leq 0, \quad (x, t) \in Q_T, \quad (16)$$

$$u(x, t) \leq 0, \quad (x, t) \in B_T \quad (17)$$

Then  $u(x, t) \leq 0$  on  $\bar{Q}_T$ .

7. Let  $a, c$  and  $c$  be as in Prob.4. Suppose that  $u(x, t)$  is continuous on  $\bar{Q}_T$  and satisfies

$$u_t - a(x, t)u_{xx} - b(x, t)u_x - c(x, t)u = 0$$

in  $Q_T$ . If  $u$  has a positive maximum , it is assumed on the boundary  $B_T$  and if  $u$  has a negative minimum , it is assumed on  $B_T$ .

8. Let  $a, b$  and  $c$  be as in Prob. 4. Suppose that  $F(x, t)$  is defined in  $Q_T$  and bounded by the constant  $N$  there. Suppose that  $u(x, t)$  is continuous on  $\bar{Q}_T$ ,  $M = \max_{B_T} |u(x, t)|$ , and  $u(x, t)$  satisfies

$$u_t - a(x, t)u_{xx} - b(x, t)u_x - c(x, t)u = F(x, t)$$

in  $Q_T$ . The  $|u(x, t)| \leq M + TN$  in  $Q_T$ .

9. Let  $a(x, t) > 0$  and  $c(x, t) \leq 0$  for all  $0 < x < L$  and  $t > 0$ ). Show that the initial boundary value problem

$$u_t - a(x, t)u_{xx} - b(x, t)u_x - c(x, t)u = F(x, t), \quad 0 < x, L, \quad t > 0, \quad (18)$$

$$u(0, t) = g(t), \quad u(L, t) = h(t), \quad t \leq 0, \quad (19)$$

$$u(x, 0) = f(x), \quad 0 \leq x \leq L \quad (20)$$

has at most one solution which is continuous on  $0 \leq x \leq L$  and  $t \geq 0$  [Notice that if there is , in fact a continuous solution, then the boundary and initial data must be compatible. That is , $g(0) = f(0)$ ,  $h(0) = f(L)$ ].

**HOMEWORK III (MATH 544): *March 7, 2012***  
**(For March 21, 2012)**

1. An initial value problem of the wave equation:

$$\begin{aligned} u_{tt} - c^2 u_{xx} &= 0, \quad t > 0, \quad 0 < x < L, \\ u(0, t) &= 0, \quad u_x(L, t) = 0, \quad t \geq 0, \\ u(x, 0) &= f(x), \quad u_t(x, 0) = g(x), \quad 0 \leq x \leq L \end{aligned}$$

where  $f(x)$  and  $g(x)$  are given functions (data) for all  $x \in (0, L)$ . Find the conditions on the data so that the above initial value problem is well posed.

2. Solve the following initial value problems.

$$x^2 u_{xx} + 3xy u_{xy} + 2y^2 u_{yy} = 0, \quad (x, y) \in D \subset \mathbb{R}^2 \quad (1)$$

with the following initial data:

(a)  $u(x, x\sqrt{x}) = f(x)$  and  $u_n(x, x\sqrt{x}) = g(x)$  where  $f(x)$  and  $g(x)$  are given functions of  $x$  and  $u_n$  is the directional derivative of  $u$  along the normal vector of the initial curve.

(b)  $u(x, x) = f(x)$  and  $u_n(x, x) = g(x)$  where  $f(x)$  and  $g(x)$  are given functions of  $x$  and  $u_n$  is the directional derivative of  $u$  along the normal vector of the initial curve.

Solution to the second Question

$$2) \quad x^2 u_{xx} + 3xy u_{xy} + 2y^2 u_{yy} = 0$$

$$\xi = \frac{x}{y}, \quad \eta = \frac{x^2}{y}$$

DE changes to  $u(x,y) = \zeta(\xi, \eta)$

$$-\eta \zeta_{\xi\eta} + \zeta_{\xi\xi} = 0$$

$$\Rightarrow \zeta(\xi, \eta) = \eta F(\xi) + G(\eta)$$

where  $F$  and  $G$  are arbitrary functions. Then

$$u(x,y) = \frac{x^2}{y} F(x/y) + G(x^2/y)$$

a) Curve is not a characteristic  $x=t, y=t^{3/2}$

$$\sqrt{t} F\left(\frac{1}{\sqrt{t}}\right) + G(\sqrt{t}) = f(t)$$

$$-(3+\frac{1}{t})F\left(\frac{1}{\sqrt{t}}\right) + 2t\left(\frac{3}{2} + \frac{1}{t}\right)F_t$$

$$-2\sqrt{t}(3+\frac{1}{t})G_t = \sqrt{1+\frac{9}{4}t} g(4)$$

$$\begin{cases} \hat{t} = \frac{1}{\sqrt{1+\frac{9}{4}t}} & (1, \frac{3}{2}\sqrt{t}) \\ \hat{n} = \frac{1}{\sqrt{1+\frac{9}{4}t}} & (-\frac{3}{2}\sqrt{t}, 1) \end{cases}$$

$$\Rightarrow F_t = \frac{\sqrt{g + 2\sqrt{t}(3+\frac{1}{t})f'}}{2t(9/4 + 1/t)} \quad "F(1/\sqrt{t})"$$

$$G_t = f' - \frac{1}{2} \frac{1}{\sqrt{t}} F - \sqrt{t} F_t \quad "G(\sqrt{t})"$$

then we obtain  $F(t)$  and  $G(t)$  in terms of data

b) Curve is a characteristic  $x=t, y=t$

$$\text{let } F(1) + G(t) = f(t)$$

$$\frac{1}{\sqrt{2}} \left[ -\frac{xy}{y^2} F - \frac{x^2}{y^2} F' - \frac{x^2}{y^2} G' \right]$$

$$t = \frac{1}{\sqrt{2}} (1, 1)$$

$$n = \frac{1}{\sqrt{2}} (-1, 1)$$

$$\frac{1}{\sqrt{2}} \left[ -\frac{xy}{y^2} F - \frac{x^2}{y^2} F' - \frac{xy}{y^2} G_y - \frac{x^2}{y^2} F' - \frac{x^3}{y^3} F_3 \right. \\ \left. - \frac{x^2}{y^2} G_y \right] = g(t) \quad \text{on } \delta$$

$$\frac{1}{\sqrt{2}} \left[ -3F(1) - 2F_3(1) - 2G_t(t) \right] = g(t)$$

$$-3F(1) - 2F_3(1) - 2G_t(t) = \sqrt{2} g(t)$$

$$\text{but } G(t) = -tF(1) + f(t)$$

$$G_t = -F(1) + f'$$

$$\Rightarrow -3F(1) - 2F_3(1) + 2F(1) - 2f' = 2g(t)$$

$$-F(1) - 2F_3(1) - 2f' = 2g(t)$$

$$\text{differentiate: } f'' = -g'$$

If data satisfied this reln we get uniquely  
many solns otherwise there exist no such