

1.(a) The image will be at $z = -a$, and will have the same magnitude, and will also point in the $+z$ direction.

(b) Since $\vec{E} = (3(\vec{p} \cdot \hat{r})\hat{r} - \vec{p})/4\pi\epsilon_0 r^3$, we will have $\vec{E}(z) = 2p_o \hat{z}/[4\pi\epsilon_0(z+a)^3]$.

(c) Since \vec{E} and \vec{p} are in the same direction, torque on the dipole will be zero.

(d) Since $\vec{F} = (\vec{p} \cdot \vec{\nabla})\vec{E}$, we have $\vec{F} = p_o (\partial/\partial z \vec{E})_{(0,0,a)} = -\hat{z}3p_o^2/(32\pi\epsilon_0 a^4)$.

The x and y components of the field are proportional to x and y respectively and do not contribute to the force at $(0,0,a)$.

2. (a) $\rho_{\text{bound}} = -\nabla \cdot \vec{P} = -\partial/\partial z \alpha z = -\alpha$.

(b) $\sigma_{\text{bound}} = \vec{P} \cdot \hat{r}|_S = P(R, \theta) \hat{z} \cdot \hat{r} = \alpha \underbrace{R \cos(\theta)}_z \cdot \underbrace{\cos(\theta)}_{\hat{z} \cdot \hat{r}}$

(c) Since $P_2(x) = (3x^2 - 1)/2$, we have $x^2 = (2P_2(x) + 1)/3$

so that $\sigma_{\text{bound}} = \alpha R \cos^2(\theta) = \alpha R [2P_2(\cos \theta) + P_0(\cos \theta)]/3$.

(d) We will have only $l = 0$ and $l = 2$ contributions:

$$V_{\text{in}}(r, \theta) = A_0 + A_2 r^2 P_2(\cos \theta) \quad V_{\text{out}}(r, \theta) = B_0/r + B_2 P_2(\cos \theta)/r^3$$

with the boundary conditions

$$V_{\text{in}}(r, \theta)|_{r=R} = V_{\text{out}}(r, \theta)|_{r=R} \quad \text{and} \quad \underbrace{\left(-\frac{\partial}{\partial r} V_{\text{out}}(r, \theta)\right)_{r=R}}_{\hat{r} \cdot \vec{E}_{\text{out}}} - \underbrace{\left(-\frac{\partial}{\partial r} V_{\text{in}}(r, \theta)\right)_{r=R}}_{\hat{r} \cdot \vec{E}_{\text{in}}} = \sigma/\epsilon_0$$

which results in

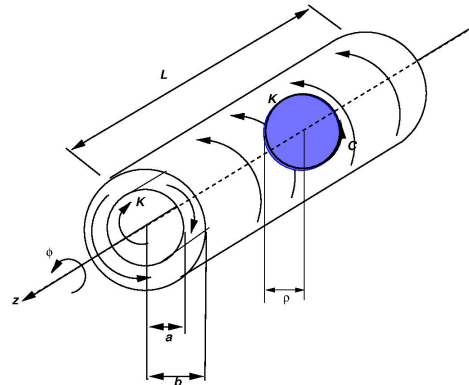
$$\begin{aligned} A_0 &= B_0/R & A_2 R^2 &= B_2/R^3 \\ B_0/R^2 &= \alpha R/3\epsilon_0 & 3B_2/R^4 + 2A_2 R &= 2\alpha R/3\epsilon_0 \end{aligned}$$

so that $B_0 = \alpha R^3/3\epsilon_0$, $A_0 = \alpha R^2/3\epsilon_0$, $B_2 = 2\alpha R^5/15\epsilon_0$ and $A_2 = 2\alpha/15\epsilon_0$

3. (a) Using superposition, and noting that the solenoidal current produces a magnetic field only inside the cylinder, we obtain the field structure $\vec{B} = \hat{z}\mu_o K$ for $a < \rho < b$ and zero elsewhere.

(b) Applying $\oint_c \vec{A} \cdot d\vec{l} = \int_S \vec{B} \cdot d\vec{S}$ to the blue surface S (bounded by the co-axial circle C with radius ρ) in the figure,

$$\oint_c \vec{A} \cdot d\vec{l} = A_\phi 2\pi\rho = \begin{cases} 0 & \text{for } \rho < a \\ B\pi(\rho^2 - a^2) & \text{for } a < \rho < b \\ B\pi(b^2 - a^2) & \text{for } \rho > b \end{cases}$$



(c) Noticing that \vec{A} is in the ϕ direction, and a function of ρ only, the only term that will give a finite contribution is $\vec{B} = \hat{z}\frac{1}{\rho}\frac{\partial}{\partial\rho}(\rho A_\phi)$. From the form of ρA_ϕ it is quite apparent that one will get zero unless $a < \rho < b$, and the differentiation will then give \vec{B} found in part (a).