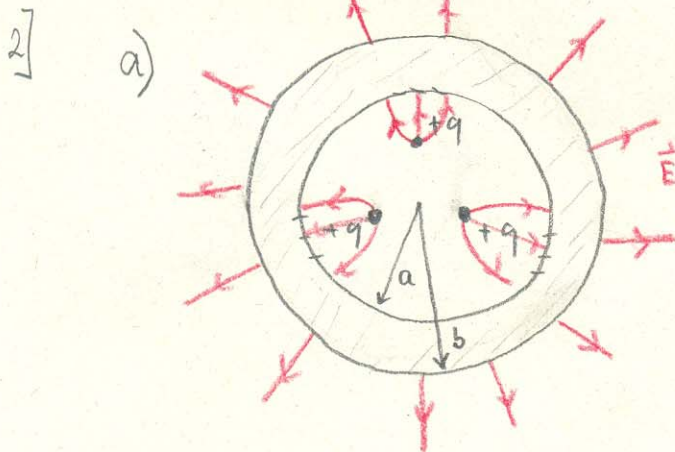


1] a) $\sigma = \epsilon_0 E = 8.85 \cdot 10^{-12} \cdot (-150) \approx -1.3 \cdot 10^{-9} \text{ C/m}^2$

b) Total charge $q = \sigma 4\pi R_{\text{Earth}}^2 = -1.3 \cdot 10^{-9} 4\pi \cdot (6.4 \cdot 10^6)^2 \approx -700 \cdot 10^3 \text{ C}$

$N = \frac{q}{e} \approx 4 \cdot 10^{24}$

c) The excess negative electric charge should be compensated by the +ve ions spread over the atmosphere. Therefore outside the atmosphere there will be (almost) no E-field. Thunderstorms are caused by the local build up of excess charge.

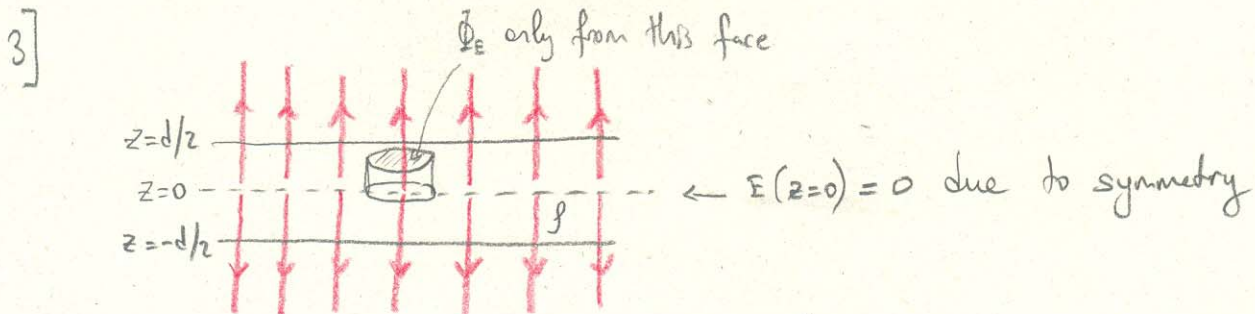


b) $V(b) = V(a) \dots$ conductor

For $r > b$, the problem is equivalent to $+3q$ charge at center

$\therefore V(b) = \frac{3q}{4\pi\epsilon_0} \frac{1}{b}$

$\Rightarrow V(a) = \frac{3q}{4\pi\epsilon_0} \frac{1}{b}$



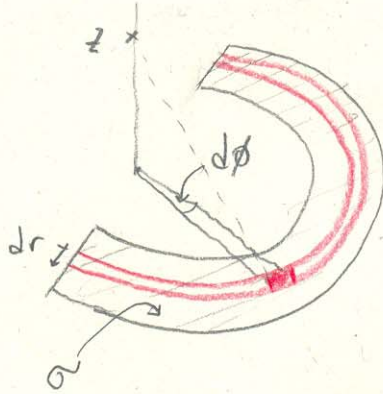
Choosing the pillbox shaped Gaussian surfaces with one face at $z=0$ plane,

we obtain
$$\oint_S \vec{E} \cdot \hat{n} da = \begin{cases} \frac{A\sigma z}{\epsilon_0}, & z \leq d/2 \\ \frac{A\sigma d}{2\epsilon_0}, & z > d/2 \end{cases}$$

$$\Rightarrow \vec{E}(z) = \begin{cases} \hat{k} \frac{\rho d}{2\epsilon_0} & , z > d/2 \\ \hat{k} \frac{\rho z}{\epsilon_0} & , |z| < d/2 \\ -\hat{k} \frac{\rho d}{2\epsilon_0} & , z < -d/2 \end{cases}$$

A]

a)



$$V(z) = \frac{1}{4\pi\epsilon_0} \int_a^b \int_0^\pi \frac{\sigma r dr d\phi}{\sqrt{z^2 + r^2}}$$

$$V(z) = \frac{\pi\sigma}{4\pi\epsilon_0} \int_a^b \frac{r dr}{\sqrt{z^2 + r^2}}$$

let $u = z^2 + r^2$; $du = 2r dr$

$$\Rightarrow V(z) = \frac{\sigma}{4\epsilon_0} \int_{z^2+a^2}^{z^2+b^2} u^{-1/2} \frac{du}{2}$$

$$\sqrt{z^2+b^2} - \sqrt{z^2+a^2}$$

$$\Rightarrow V(z) = \frac{\sigma}{4\epsilon_0} \left[\sqrt{z^2+b^2} - \sqrt{z^2+a^2} \right]$$

b) $E_z(z) = - \frac{\partial V(z)}{\partial z}$

$$E_z(z) = \frac{\sigma z}{4\epsilon_0} \left[\frac{1}{\sqrt{z^2+a^2}} - \frac{1}{\sqrt{z^2+b^2}} \right]$$

Note that, \vec{E} -field has other components as well.