

**E33-9** For a single long straight wire,  $B = \mu_0 i / 2\pi d$  but we need a factor of “2” since there are two wires, then  $i = \pi dB / \mu_0$ . Finally

$$i = \frac{\pi dB}{\mu_0} = \frac{\pi(0.0405 \text{ m})(296, \mu\text{T})}{(4\pi \times 10^{-7} \text{ N/A}^2)} = 30 \text{ A}$$

**E33-15** We imagine the ribbon conductor to be a collection of thin wires, each of thickness  $dx$  and carrying a current  $di$ .  $di$  and  $dx$  are related by  $di/dx = i/w$ . The contribution of one of these thin wires to the magnetic field at  $P$  is  $dB = \mu_0 di/2\pi x$ , where  $x$  is the distance from this thin wire to the point  $P$ . We want to change variables to  $x$  and integrate, so

$$B = \int dB = \int \frac{\mu_0 i dx}{2\pi w x} = \frac{\mu_0 i}{2\pi w} \int \frac{dx}{x}.$$

The limits of integration are from  $d$  to  $d + w$ , so

$$B = \frac{\mu_0 i}{2\pi w} \ln \left( \frac{d + w}{d} \right).$$

**P33-14** (a) According to Eq. 33-34, the magnetic field inside the wire *without a hole* has magnitude  $B = \mu_0 i r / 2\pi R^2 = \mu_0 j r / 2$  and is directed radially. If we superimpose a second current to create the hole, the additional field at the center of the hole is zero, so  $B = \mu_0 j b / 2$ . But the current in the remaining wire is

$$i = jA = j\pi(R^2 - a^2),$$

so

$$B = \frac{\mu_0 i b}{2\pi(R^2 - a^2)}.$$

**P34-8** (a)  $d\Phi_B/dt = B dA/dt$ , but  $dA/dt$  is  $\Delta A/\Delta t$ , where  $\Delta A$  is the area swept out during one rotation and  $\Delta t = 1/f$ . But the area swept out is  $\pi R^2$ , so

$$|\mathcal{E}| = \frac{d\Phi_B}{dt} = \pi f B R^2.$$

(b) If the output current is  $i$  then the power is  $P = \mathcal{E}i$ . But  $P = \tau\omega = \tau 2\pi f$ , so

$$\tau = \frac{P}{2\pi f} = i B R^2 / 2.$$

**P34-13** Assume that  $E$  does vary as the picture implies. Then the line integral along the path shown *must* be nonzero, since  $\vec{\mathbf{E}} \cdot d\vec{\mathbf{l}}$  on the right is not zero, while it is along the three other sides. Hence  $\oint \vec{\mathbf{E}} \cdot d\vec{\mathbf{l}}$  is non zero, implying a change in the magnetic flux through the dotted path. But it doesn't, so  $\vec{\mathbf{E}}$  cannot have an abrupt change.

**E36-8** In each case apply  $\mathcal{E} = L\Delta i/\Delta t$ .

(a)  $\mathcal{E} = (4.6 \text{ H})(7 \text{ A})/(2 \times 10^{-3} \text{ s}) = 1.6 \times 10^4 \text{ V}$ .

(b)  $\mathcal{E} = (4.6 \text{ H})(2 \text{ A})/(3 \times 10^{-3} \text{ s}) = 3.1 \times 10^3 \text{ V}$ .

(c)  $\mathcal{E} = (4.6 \text{ H})(5 \text{ A})/(1 \times 10^{-3} \text{ s}) = 2.3 \times 10^4 \text{ V}$ .

**P36-2** (a) Since  $ni$  is the net current per unit length and in this case  $i/W$ , we can simply write  $B = \mu_0 i/W$ .

(b) There is only one loop of wire, so

$$L = \phi_B/i = BA/i = \mu_0 i \pi R^2 / Wi = \mu_0 \pi R^2 / W.$$

**P36-3** Choose the  $y$  axis so that it is parallel to the wires and directly between them. The field in the plane between the wires is

$$B = \frac{\mu_0 i}{2\pi} \left( \frac{1}{d/2 + x} + \frac{1}{d/2 - x} \right).$$

The flux per length  $l$  of the wires is

$$\begin{aligned} \Phi_B &= l \int_{-d/2+a}^{d/2-a} B dx = l \frac{\mu_0 i}{2\pi} \int_{-d/2+a}^{d/2-a} \left( \frac{1}{d/2 + x} + \frac{1}{d/2 - x} \right) dx, \\ &= 2l \frac{\mu_0 i}{2\pi} \int_{-d/2+a}^{d/2-a} \left( \frac{1}{d/2 + x} \right) dx, \\ &= 2l \frac{\mu_0 i}{2\pi} \ln \frac{d-a}{a}. \end{aligned}$$

The inductance is then

$$L = \frac{\phi_B}{i} = \frac{\mu_0 l}{\pi} \ln \frac{d-a}{a}.$$